

Dr. Mohamed Bakr, EE3FK4, 2008

Note Title

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# Lecture 11

From Sections 10.3 of Textbook

Solve 10.2 - 10.10

## Wave Equation Using Harmonic Fields

$$\nabla \times \underline{\underline{E}} = -j\omega \underline{\underline{H}}, \quad \nabla \times \underline{\underline{H}} = j\omega \underline{\underline{E}}$$

Take the  $\nabla \times$  of the first equation to

$$\text{have } \nabla \times \nabla \times \underline{\underline{E}} = -j\omega \nabla \times \underline{\underline{H}}$$

$$\underline{\underline{j\omega \epsilon E}}$$

$$\Rightarrow \nabla \times \nabla \times \underline{\underline{E}} = \omega^2 \mu \epsilon \underline{\underline{E}}$$

$$\nabla(\nabla \cdot \underline{\underline{E}}) - \nabla^2 \underline{\underline{E}} = \omega^2 \mu \epsilon \underline{\underline{E}}$$

$$\Rightarrow \text{For no sources } \nabla \cdot \underline{\underline{E}} = 0$$

## Wave Equation (Cont'd)

$$\Rightarrow \nabla^2 \underline{E} + \beta^2 \underline{E} = 0, \quad \beta = \omega \sqrt{\mu \epsilon} = \frac{\omega}{c_p}$$

For x Component  $\nabla^2 E_x + \beta^2 E_x = 0$

$$\Rightarrow \frac{\partial^2 E_x}{\partial x^2} + \frac{\partial^2 E_x}{\partial y^2} + \frac{\partial^2 E_x}{\partial z^2} + \beta^2 E_x = 0$$

\* Now if the field has only an x Component and it is only a function of  $z$ , we

$$\text{Have } \frac{\partial^2 E_x}{\partial z^2} + \beta^2 E_x = 0$$

## Wave Equation (Cont'd)

$$E_x = E_{x_0}^+ e^{-j\beta z} + E_{x_0}^- e^{j\beta z}$$

+ve z wave                      -ve z wave

\* Since  $E_x$  is known,  $H_y$  can be obtained

Using Maxwell's equations

$$\nabla \times \underline{E} = -j\omega \underline{H}$$

$$\left( \frac{\partial E_x}{\partial z} \right) \underline{a}_y = -j\omega \underline{H} = -j\omega H_y \underline{a}_y$$

$$H_y = -\frac{1}{j\omega} (-j\beta E_{x_0}^+ e^{-j\beta z} + j\beta E_{x_0}^- e^{j\beta z})$$

Wave Equation (Cont'd)

$$\vec{H}_y = \frac{\beta}{\omega\mu} E_{x_0}^+ e^{-jkz} - \frac{\beta}{\omega\mu} E_{x_0}^- e^{jkz}$$

$$\vec{H}_y = H_{y_0}^+ e^{-jkz} + H_{y_0}^- e^{jkz}$$

$$\frac{E_{x_0}^+}{H_{y_0}^+} = \frac{\omega\mu}{\beta} = \frac{\omega\mu}{\frac{\omega}{v}} = v * \frac{1}{\mu\epsilon} = \sqrt{\frac{\mu}{\epsilon}} \approx$$

$$\text{also } - \frac{E_{x_0}^-}{H_{y_0}^-} = \frac{\omega\mu}{\beta} = \sqrt{\frac{\mu}{\epsilon}}$$

$$\text{for air } \mu_0 = \sqrt{\frac{\mu_0}{\epsilon_0}} = 377 \Omega$$

## Time Domain Expression

$$\vec{E}_x(z) = E_{x_0}^+ e^{-j\beta z} + E_{x_0}^- e^{j\beta z}$$

$$\vec{E}_x(z) = |E_{x_0}^+| e^{-j\beta z + j\phi} + |E_{x_0}^-| e^{j\beta z + j\phi}$$

$$E_x(z, t) = |E_{x_0}^+| \cos(\omega t - \beta z + \phi) + |E_{x_0}^-| \cos(\omega t + \beta z + \phi)$$

+z wave

-z wave

+z wave

-z wave



## Propagation in lossy media

\* This Derivation is much easier to do in the frequency domain

$$\nabla \times \underline{\underline{E}} = -j\omega \underline{\underline{H}}, \quad \nabla \times \underline{\underline{H}} = \sigma \underline{\underline{E}} + j\omega \underline{\underline{\epsilon}} \underline{\underline{E}}$$

$$\Rightarrow \nabla \times \underline{\underline{H}} = j\omega \underline{\underline{\epsilon}} \left(1 - j \frac{\sigma}{\omega \underline{\underline{\epsilon}}}\right) \underline{\underline{E}}$$

$$\nabla \times \underline{\underline{H}} = j\omega \left(\underline{\underline{\epsilon}} - j \frac{\sigma}{\omega}\right) \underline{\underline{E}} = j\omega \left(\underline{\underline{\epsilon}} - j \underline{\underline{\epsilon}}''\right) \underline{\underline{E}}$$

\* Losses can be modelled using complex permittivity!

## Propagation in Lossy Media (Cont'd)

$$K = \omega \sqrt{\mu \epsilon} = \omega \sqrt{\mu (\epsilon' - j\epsilon'')}$$

$$K = \omega \sqrt{\mu \epsilon'} \sqrt{1 - j\frac{\epsilon''}{\epsilon'}}$$

So, we can obtain

$$\alpha = \text{Re}(jK) = \omega \sqrt{\frac{\mu \epsilon'}{2}} \left( \sqrt{1 + \left(\frac{\epsilon''}{\epsilon'}\right)^2} - 1 \right)^{1/2}$$

$$\beta = \text{Im}(jK) = \omega \sqrt{\frac{\mu \epsilon'}{2}} \left( \sqrt{1 + \left(\frac{\epsilon''}{\epsilon'}\right)^2} + 1 \right)^{1/2}$$



## Propagation in a Lossy Medium (cont'd)

$$\underline{E}^z = E_x^+ e^{-jkz} = E_x^+ e^{-(\alpha + j\beta)z} e^{-\alpha z}$$

$$E(z, t) = \operatorname{Re} \left( E_x^+ e^{-(\alpha + j\beta)z} e^{j\omega t} \right) e^{-\alpha z}$$

$$E(z, t) = \operatorname{Re} \left( |E_x^+| e^{-\alpha z} e^{j(\omega t - \beta z + \phi^+)} \right) e^{-\alpha z}$$

$$E(z, t) = |E_x^+| e^{-\alpha z} \cos(\omega t - \beta z + \phi^+)$$

Wave decays as it travels in the  $z$  direction!

# Lossy Medium

