

Lecture 10: Electrostatics

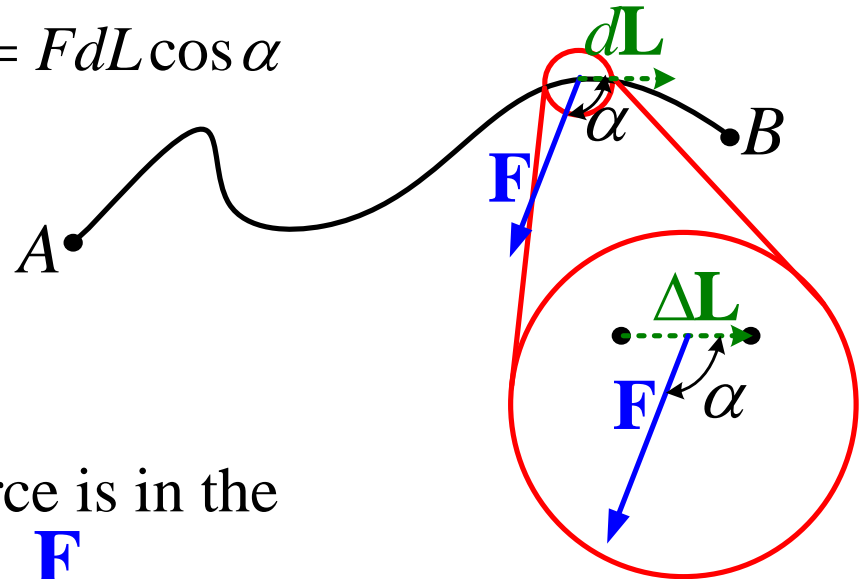
Electric potential, relations between E and V ,
Chapter 4, pages 137-146

Work

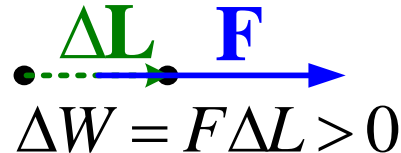
differential work $dW = \mathbf{F} \cdot d\mathbf{L} = FdL \cos \alpha$

work done from A to B

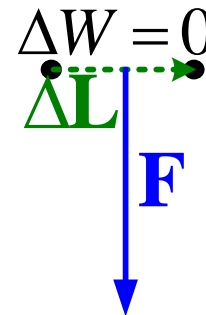
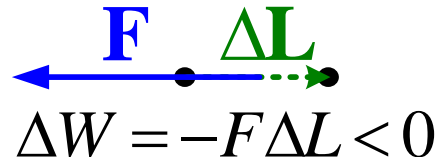
$$W = \int_A^B \mathbf{F} \cdot d\mathbf{L}, \text{ J}$$



work is positive if the applied force is in the direction of displacement

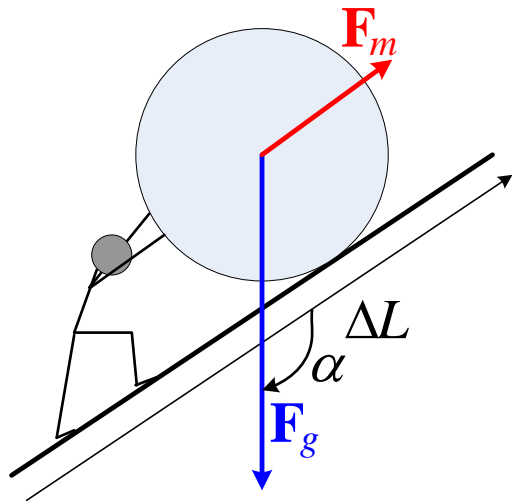


work is negative if the applied force is opposite to the direction of displacement



work is zero if force is orthogonal

Gravity Example (uniform motion)



$W_m = F_m \Delta L > 0$, man does work
(expends energy)

$W_g = F_g \Delta L \cos \alpha < 0$, rock gains energy

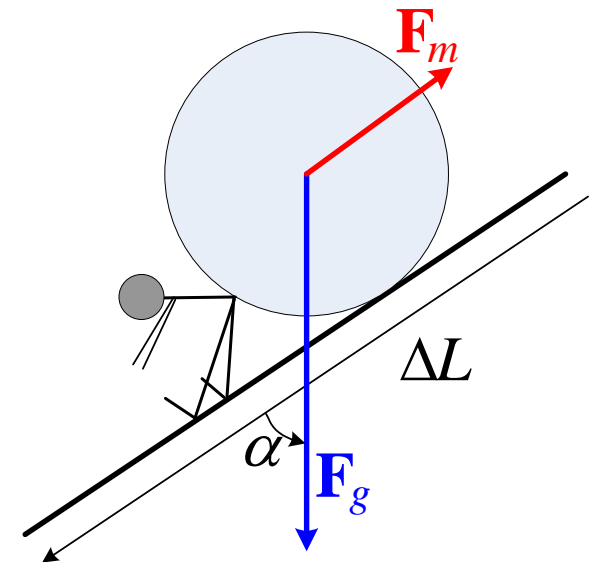
$$W_m = -W_g$$

$$F_m = -F_g \cos \alpha$$

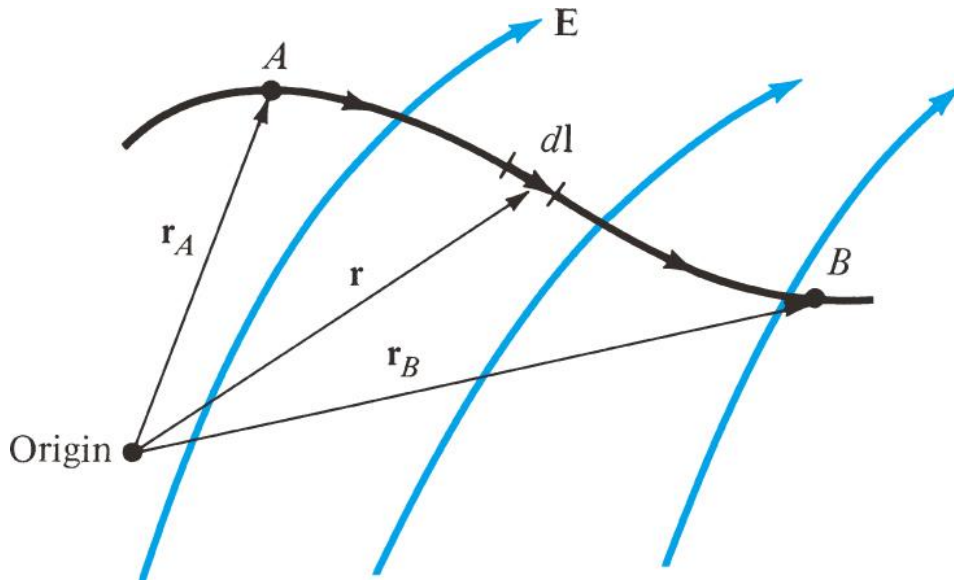
$W_g = F_g \Delta L \cos \alpha > 0$, gravity does work
(rock loses energy)

$$W_m = -F_m \Delta L < 0$$

man provides enough friction to prevent
acceleration, gains energy (heat)



Work Against Electric Field



work done against electric field to move a charge Q from point A to point B is

$$W = -Q \int_A^B \mathbf{E} \cdot d\mathbf{L} \quad \text{why -ve sign?}$$

Electrostatic Potential Difference

potential difference V_{AB} is the work done by the electric field in moving a unit test charge from point B to point A

$$V_{AB} = V_B - V_A = \int_B^A \mathbf{E} \cdot d\mathbf{L}, \text{ V}$$

OR: potential difference V_{AB} is the work done by external forces in moving a unit test charge from point A to point B

$$V_{AB} = V_B - V_A = -\int_A^B \mathbf{E} \cdot d\mathbf{L}, \text{ V}$$

Electrostatic Potential

the potential at a point A is the work done by the electric field in moving a unit test charge from point A to a reference point A_0 where the potential is assumed equal to zero

OR: the potential at a point A is the work done by an external force in moving a unit test charge from a reference point A_0 , where the potential is assumed equal to zero, to the point A .

$$V_A = \int_A^{A_0} \mathbf{E} \cdot d\mathbf{L} = - \int_{A_0}^A \mathbf{E} \cdot d\mathbf{L}, \text{ V}$$

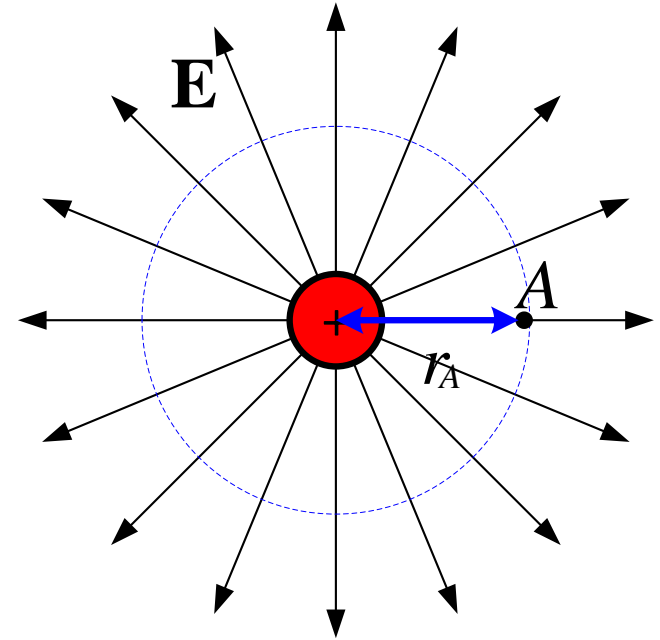
the electrostatic potential is the potential energy of a unit test charge

Potential of a Point Charge

assume potential is zero at infinity

integrate along any radial line

$$V_A = \int_A^{\infty} \mathbf{E} \cdot d\mathbf{L} = \int_{r_A}^{\infty} \frac{Q}{4\pi\epsilon r^2} \mathbf{a}_R \cdot \mathbf{a}_R dr$$
$$\Rightarrow V_A = -\frac{Q}{4\pi\epsilon} \frac{1}{r} \Big|_{r_A}^{\infty} \Rightarrow V_A = \frac{Q}{4\pi\epsilon} \frac{1}{r_A}$$



the potential is the same for all points on the sphere of radius r_A

the sphere is an equipotential surface

notice the $1/r$ dependence on distance!

Potential (Cont'd)

if the points charge is at \mathbf{r}' , the potential at a point \mathbf{r} is given by

$$V(\mathbf{r}) = \frac{Q}{4\pi\epsilon |\mathbf{r} - \mathbf{r}'|}$$

for a number of discrete charges, we have

$$V(\mathbf{r}) = \sum_{i=1}^N \frac{Q_i}{4\pi\epsilon |\mathbf{r} - \mathbf{r}'_i|}$$

remember that all these potentials are relative to a point in ∞ !

Potential (Cont'd)

for continuous charge distributions, we have

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon} \int \frac{\rho_L(\mathbf{r}') dl'}{4\pi\epsilon |\mathbf{r} - \mathbf{r}'|} \quad \text{line charge}$$

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon} \int \frac{\rho_s(\mathbf{r}') ds'}{4\pi\epsilon |\mathbf{r} - \mathbf{r}'|} \quad \text{surface charge}$$

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon} \int \frac{\rho_v(\mathbf{r}') dv'}{4\pi\epsilon |\mathbf{r} - \mathbf{r}'|} \quad \text{volume charge}$$

Potential of a Line Charge

$$V_{AB} = V_B - V_A = -\int_A^B \mathbf{E} \cdot d\mathbf{L} = -\int_{\rho_A}^{\rho_B} \frac{\rho_L}{2\pi\rho} \mathbf{a}_\rho \cdot \mathbf{a}_\rho d\rho$$

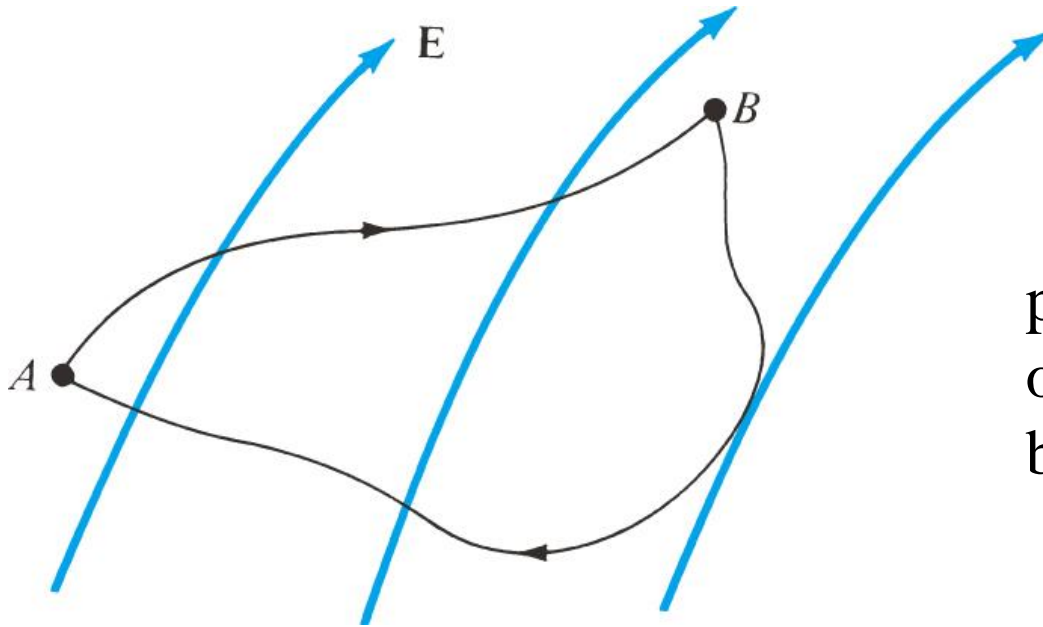
$$\Rightarrow V_{AB} = \frac{\rho_l}{2\pi\epsilon} \cdot \ln\left(\frac{\rho_A}{\rho_B}\right)$$

Can we choose ρ_B or ρ_A at infinity or at zero?

if we choose $\rho_B=1.0$ m, and $V(\rho_B)=0$, we get

$$\Rightarrow V_{AB} = -V_A = \frac{\rho_l}{2\pi\epsilon} \cdot \ln(\rho_A) \Rightarrow V_A = \frac{\rho_l}{2\pi\epsilon} \cdot \ln\left(\frac{1}{\rho_A}\right)$$

Relation between \mathbf{E} and V



potential does not depend
on contour of integration
but on position of points

$$V_{AB} = -\int_{C1} \mathbf{E} \cdot d\mathbf{L} = -\int_{C2} \mathbf{E} \cdot d\mathbf{L} \Rightarrow \oint \mathbf{E} \cdot d\mathbf{L} = 0 \quad \text{ESF is conservative!}$$

using Stoke's Theorem, $\oint \mathbf{E} \cdot d\mathbf{L} = \iint_S (\nabla \times \mathbf{E}) \cdot d\mathbf{s} = 0 \Rightarrow (\nabla \times \mathbf{E}) = \mathbf{0}$

it follows that $\mathbf{E} = -\nabla V$ (why?)

Relation between \mathbf{E} and V (Cont'd)

the differential potential difference over a differential length $d\mathbf{L}$ is given by

$$dV = -\mathbf{E} \cdot d\mathbf{L} = -E_x dx - E_y dy - E_z dz$$

using first-order Taylor expansion, we have

$$dV = \frac{\partial V}{\partial x} dx + \frac{\partial V}{\partial y} dy + \frac{\partial V}{\partial z} dz$$

equating, we have

$$E_x = -\frac{\partial V}{\partial x}, \quad E_y = -\frac{\partial V}{\partial y}, \quad E_z = -\frac{\partial V}{\partial z}$$