

EE750
Advanced Engineering Electromagnetics
Lecture 3

Special Cases in Maxwell's Equations.

- With no sources, Maxwell's Equations are written as

$$(\nabla \times \mathbf{E}) = -\frac{\partial \mathbf{B}}{\partial t}, \quad (\nabla \times \mathbf{H}) = \mathbf{J}_c + \frac{\partial \mathbf{D}}{\partial t}$$

- In Cartesian coordinates we have

$$\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = -\frac{\partial B_x}{\partial t}$$

$$\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} = -\frac{\partial B_y}{\partial t}$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = -\frac{\partial B_z}{\partial t}$$

$$\frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} = J_x + \frac{\partial D_x}{\partial t}$$

$$\frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} = J_y + \frac{\partial D_y}{\partial t}$$

$$\frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} = J_z + \frac{\partial D_z}{\partial t}$$

Special Cases (Cont'd)

- For a 1D case, the fields depend only on one coordinate (e.g. a uniform plane wave traveling in the x direction)
- For example, if $\partial/\partial y = \partial/\partial z = 0$

$$-\frac{\partial E_z}{\partial x} = -\frac{\partial B_y}{\partial t} \quad , \quad \frac{\partial H_y}{\partial x} = J_z + \frac{\partial D_z}{\partial t} \quad (\text{propagation in } -\text{ve } x)$$

or

$$\frac{\partial E_y}{\partial x} = -\frac{\partial B_z}{\partial t} \quad , \quad -\frac{\partial H_z}{\partial x} = J_y + \frac{\partial D_y}{\partial t} \quad (\text{propagation in } +\text{ve } x)$$

Notice that these two systems are decoupled

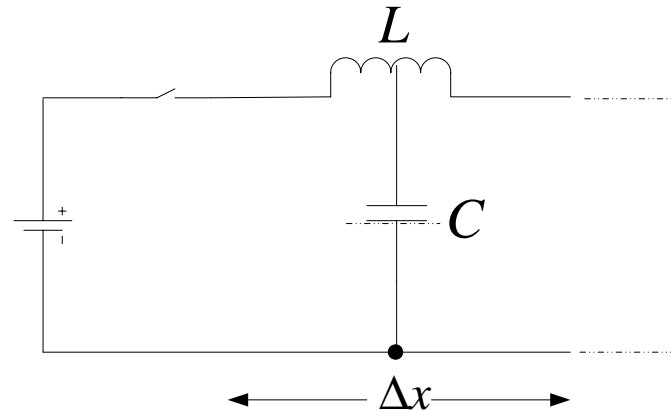
Special Cases (Cont'd)

- For the 2D problems, the fields depend only on two coordinates (e.g. TE₁₀ mode in waveguides)
- For example, if $\partial/\partial y=0$

| TE _y | | TM _y |
|--|--|---|
| $-\frac{\partial H_y}{\partial z} = J_x + \frac{\partial D_x}{\partial t}$ | | $-\frac{\partial E_y}{\partial z} = -\frac{\partial B_x}{\partial t}$ |
| $\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} = -\frac{\partial B_y}{\partial t}$ | | $\frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} = J_y + \frac{\partial D_y}{\partial t}$ |
| $\frac{\partial H_y}{\partial x} = J_z + \frac{\partial D_z}{\partial t}$ | | $\frac{\partial E_y}{\partial x} = -\frac{\partial B_z}{\partial t}$ |

Notice that the two systems are decoupled

Transient Response of a Transmission Line



- C_d capacitance per unit length F/m

L_d inductance per unit length H/m

u velocity of propagation

- After time Δt , the disturbance traveled a distance Δx .
A charge $\Delta Q = C_d \Delta x V_s$ been transferred. It follows that

$$i = \Delta Q / \Delta t \quad \Longrightarrow \quad i = C_d V_s (\Delta x / \Delta t) = C_d V_s u$$

Transient Response of a TL (Cont'd)

- The flow of current establishes a flux Φ associated with line inductance through $\Phi = L_d \Delta x i = L_d \Delta x C_d V_s u$
- Using Faraday's law we have

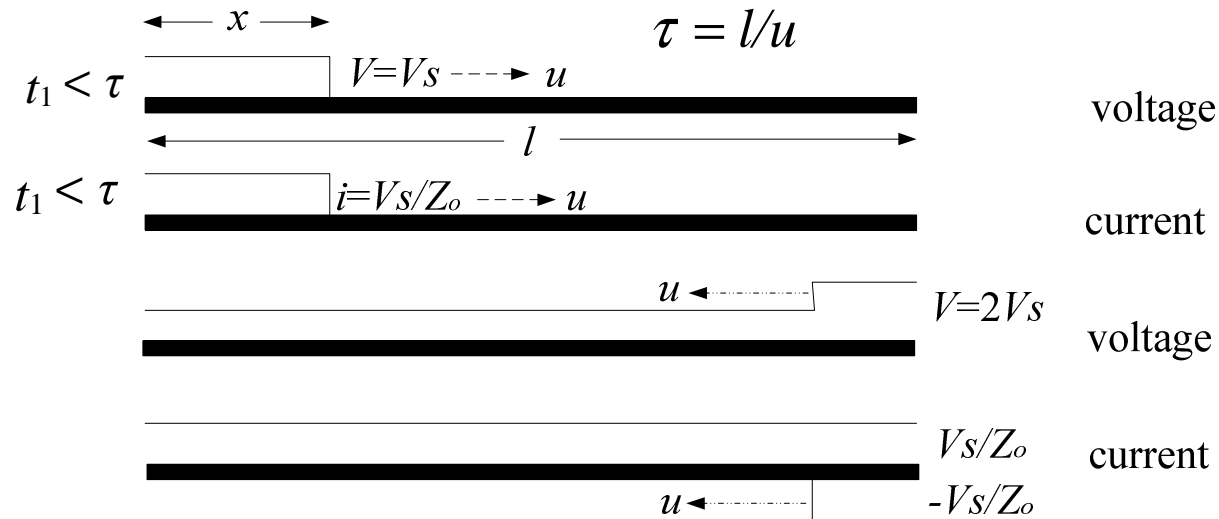
$$V_s = \Delta \Phi / \Delta t = L_d C_d V_s u^2 \implies L_d C_d V_s u^2 = 1.0$$

$$u = \frac{1}{\sqrt{C_d L_d}} \text{ m/s}$$

- It follows that $i = C_d V_s u = V_s \sqrt{\frac{C_d}{L_d}}$

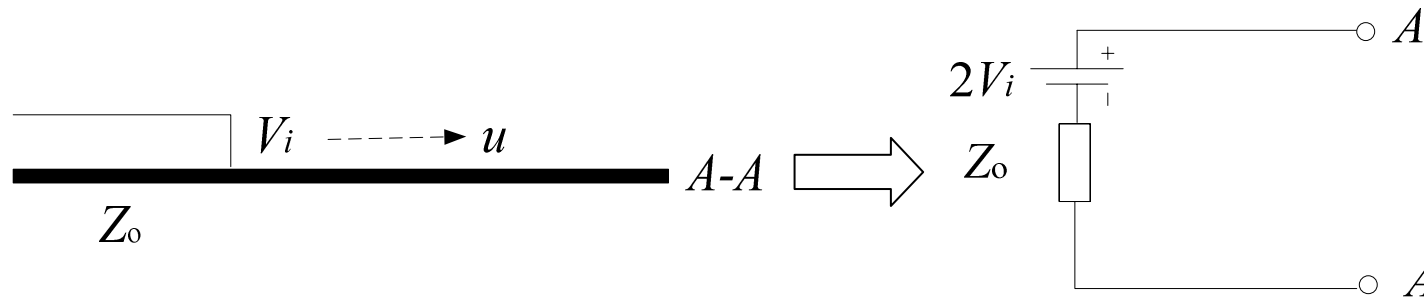
$$Z_o = \frac{\text{incident voltage}}{\text{incident current}} = \frac{V_s}{i} = \sqrt{\frac{L_d}{C_d}} \Omega$$

Open-Ended Transmission Line



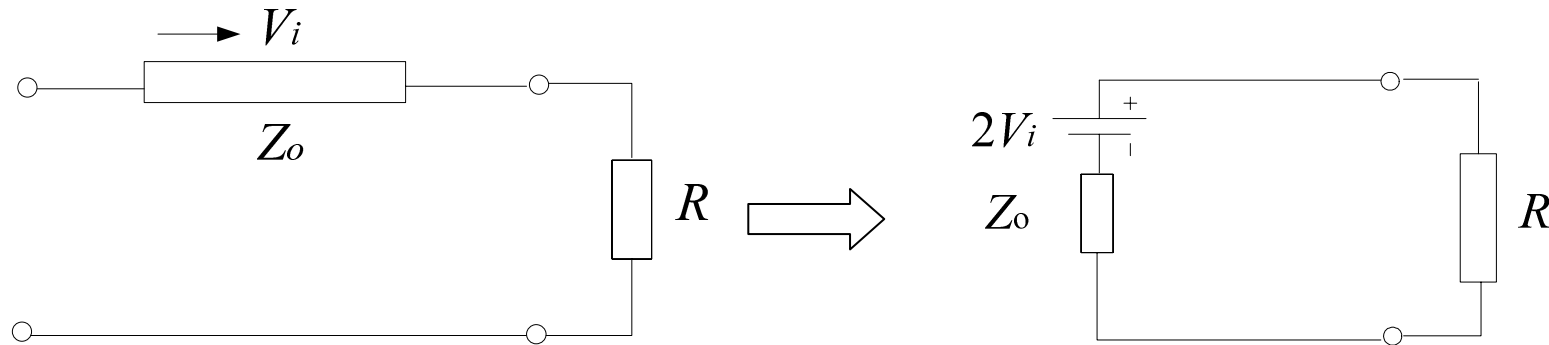
- After time τ , the incident voltage and current waves reach the open end. A reflected current wave of amplitude $I^r = -V_s/Z_o$ is initiated to make the total current at the open end zero
- It follows that for an open-ended transmission line we have $V^i = V^s$, $I^i = V_s/Z_o$ and $V^r = V^s$, $I^r = -V_s/Z_o$.

Thevenin Equivalent of a Transmission Line



- The open circuit voltage is $2V^i$ (open-ended line)
- The equivalent resistance is Z_o
- This Thevenin equivalent is valid only for limited time period where no reflections take place on the line (same V^i)

Example on Utilizing Thevenin's Equivalent of a TL

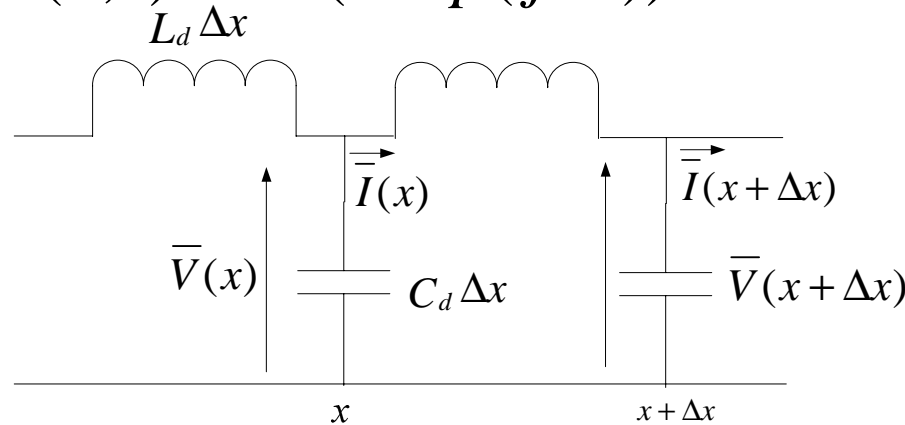


- Using Kirchoff's voltage law we have $V = 2V^i \frac{R}{R + Z_o}$
- But $V = V^i + V^r \implies V^r = V - V^i$

$$V^r = 2V^i \frac{R}{R + Z_o} - V^i \implies V^r = \left(\frac{R - Z_o}{R + Z_o} \right) V^i = \Gamma V^i$$
- For open circuit termination $R = \infty$, $\Gamma = 1$, $V^r = V^i$
- For short circuit termination $R = 0$, $\Gamma = -1$, $V^r = -V^i$

Sinusoidal Steady-State Response of a TL

- Remember that for a steady-state sinusoidal analysis we need only the amplitude and phase at each point on the line $V(x,t) = \text{Re}(\bar{V} \exp(j\omega t))$



- Applying KVL, we get $\bar{V}(x) - \bar{V}(x + \Delta x) = j\omega L_d \Delta x \bar{I}(x)$

Leading to
$$-\frac{\partial \bar{V}(x)}{\partial x} = j\omega L_d \bar{I}(x)$$

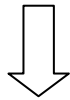
- Applying KCL we get $\bar{I}(x) - \bar{I}(x + \Delta x) = j\omega C_d \Delta x \bar{V}(x)$

Leading to
$$-\frac{\partial \bar{I}(x)}{\partial x} = j\omega C_d \bar{V}(x)$$

Sinusoidal Steady State of a TL (Cont'd)

- Differentiating to eliminate the current term we get

$$\frac{\partial^2 \bar{V}(x)}{\partial x^2} = -\omega^2 L_d C_d \bar{V}(x) = -\beta^2 \bar{V}(x)$$



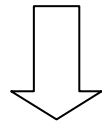
$$\bar{V}(x) = V^i \exp(-j\beta x) + V^r \exp(j\beta x)$$

- Similarly $\bar{I}(x) = I^i \exp(-j\beta x) + I^r \exp(j\beta x)$
- Using the telegrapher's differential equations we can show that $I^i = V^i / Z_o$ and $I^r = -V^r / Z_o$.
- It follows that $\bar{V}(x) = V^i \exp(-j\beta x) + V^r \exp(j\beta x)$
and $\bar{I}(x) = (V^i / Z_o) \exp(-j\beta x) - (V^r / Z_o) \exp(j\beta x)$

ABCD Matrix of a TL Section

- At $x=0$, we have $\bar{V}(0) = V^i + V^r$ and $\bar{I}(0) = (V^i / Z_o) - (V^r / Z_o)$
- Expressing the solutions in terms of the amplitudes at $x=0$, we get

$$\begin{bmatrix} \bar{V}(x) \\ \bar{I}(x) \end{bmatrix} = \begin{bmatrix} \cos(\beta x) & -j Z_o \sin(\beta x) \\ -j \sin(\beta x) / Z_o & \cos(\beta x) \end{bmatrix} \begin{bmatrix} \bar{V}(0) \\ \bar{I}(0) \end{bmatrix}$$

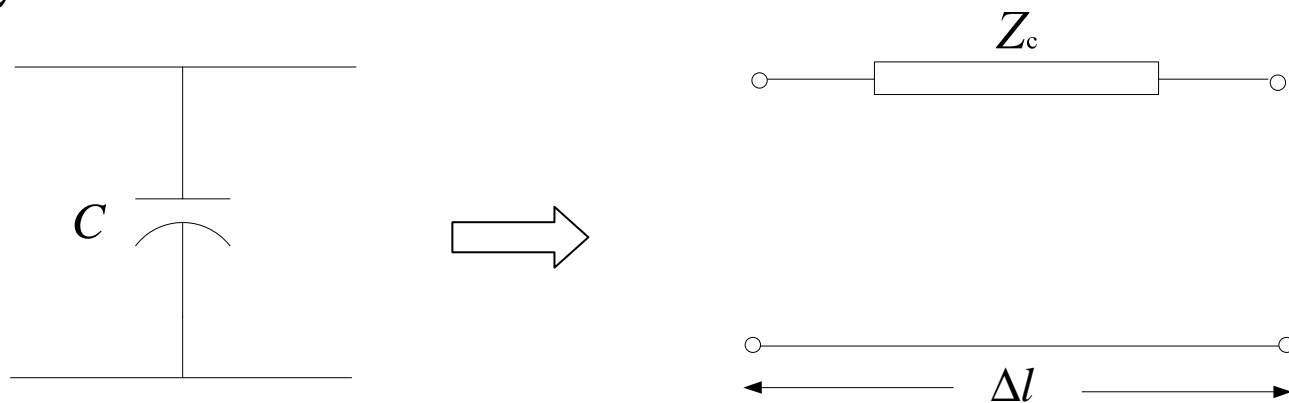


$$\begin{bmatrix} \bar{V}(0) \\ \bar{I}(0) \end{bmatrix} = \begin{bmatrix} \cos(\beta x) & j Z_o \sin(\beta x) \\ j \sin(\beta x) / Z_o & \cos(\beta x) \end{bmatrix} \begin{bmatrix} \bar{V}(x) \\ \bar{I}(x) \end{bmatrix}$$

ABCD matrix of a TL section with length x

Discrete Models of Lumped Elements

- The lumped element is replaced by a section of a transmission line
- Using these models the voltages and currents are obtained only at discrete instants of time



- $C_d \Delta l = C$, C_d is the capacitance per unit length
- $u = \Delta l / \Delta t = 1 / \sqrt{L_d C_d} \implies L_d = \left(\frac{\Delta t}{\Delta l} \right)^2 \frac{1}{C_d}$

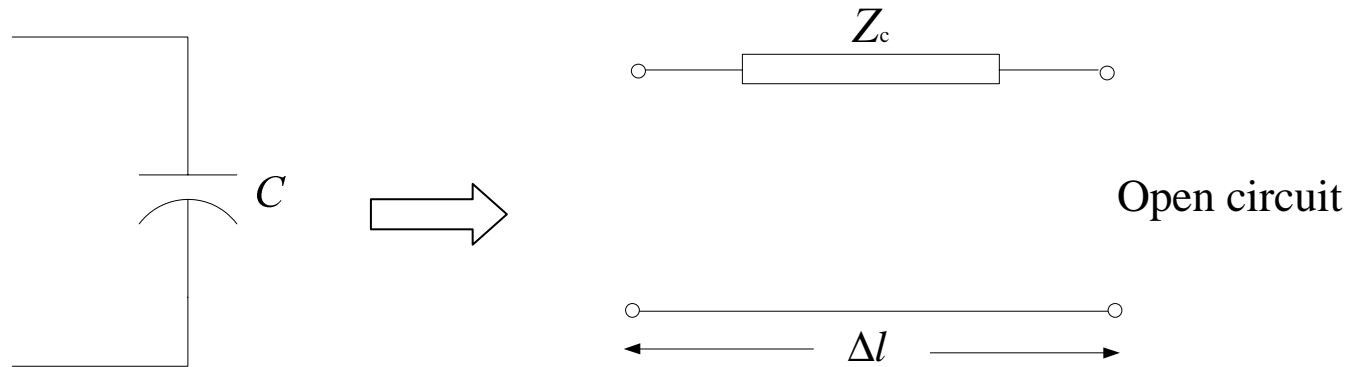
Discrete Models of Lumped Elements (Cont'd)

- $Z_c = \sqrt{\frac{L_d}{C_d}} = \frac{\Delta t}{C_d \Delta l} = \frac{\Delta t}{C}$
- $L_e =$ equivalent (parasitic) inductance $= L_d \Delta l = \frac{(\Delta t)^2}{C}$
- Δt is selected small enough such that the parasitic inductance is small

Modeling Steps:

1. Select Δt , Δl
2. evaluate C_d , L_d and Z_c
3. Assume certain incident voltage at the k th time step
4. Use Thevenin's equivalent to get the reflected voltages
5. Obtain incident impulses at the $(k+1)$ time step

Discrete Models of Lumped Elements (Cont'd)



- Using a similar derivation $C_d \Delta l = C$

$u = \Delta l / (\Delta t / 2) = 1 / \sqrt{L_d C_d}$, notice that Δt is the round-trip time

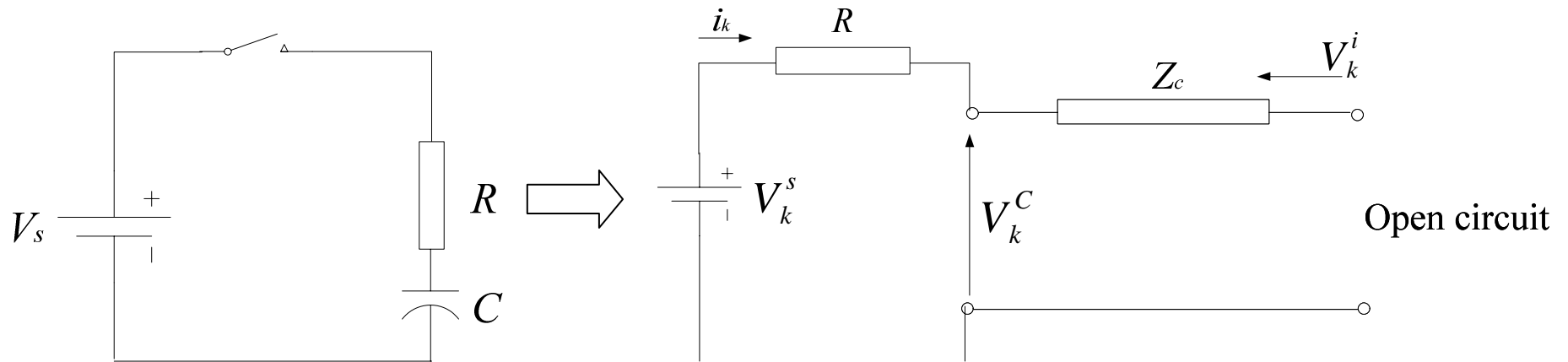
$$L_d = \left(\frac{\Delta t}{2\Delta l} \right)^2 \frac{1}{C_d} = \frac{\Delta t^2}{4C\Delta l}$$

$$Z_c = \sqrt{\frac{L_d}{C_d}} = \frac{\Delta t}{2C}$$

$$L_e = L_d \Delta l = \frac{(\Delta t)^2}{4C} = \text{error inductance}$$

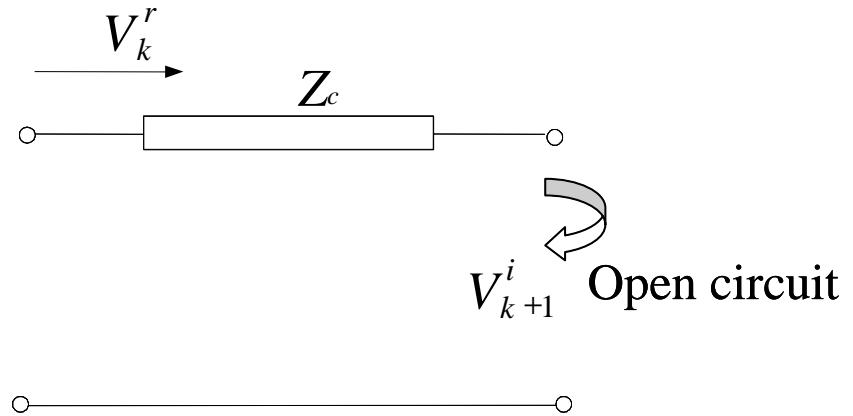
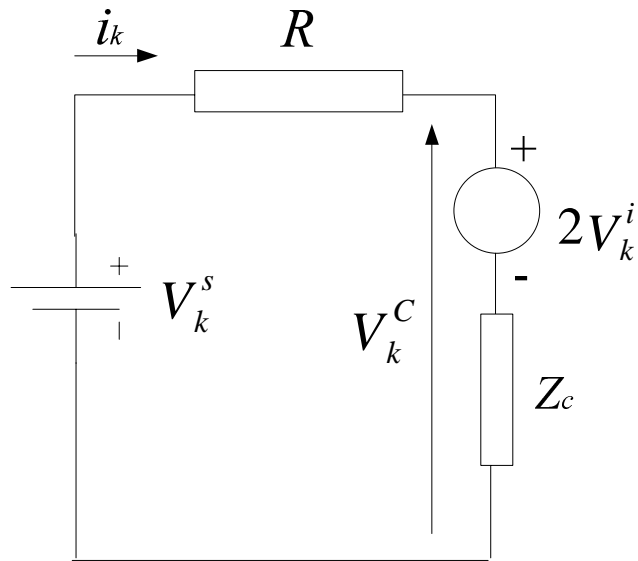
Example

- Evaluate the transient response of the shown circuit

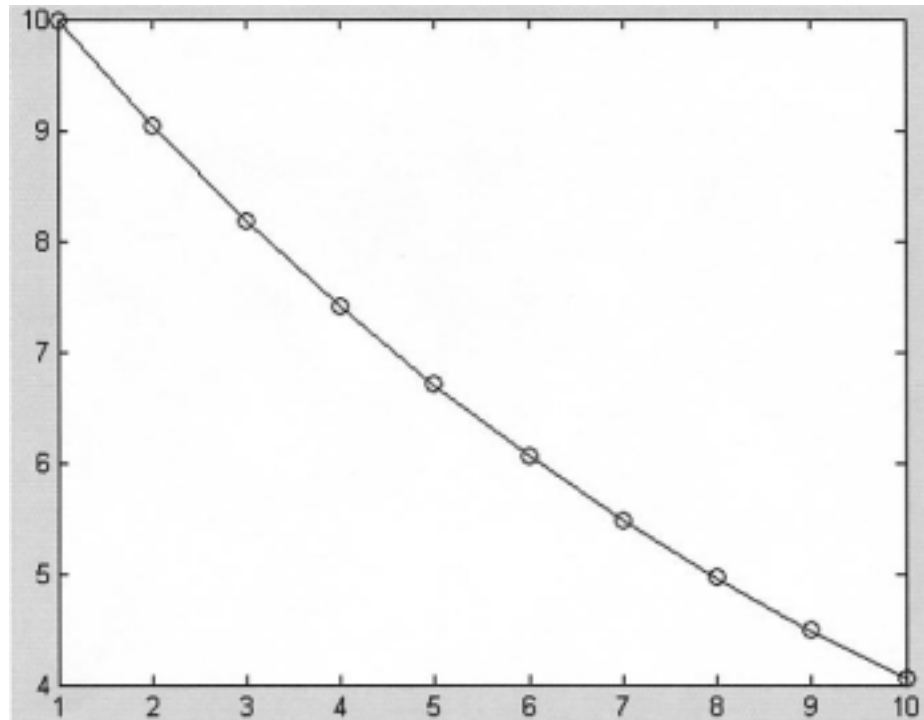


- Evaluate the current $i_k = (V_k^S - 2V_k^i) / (Z_c + R)$
- Evaluate the capacitor voltage $V_k^C = 2V_k^i + Z_c i_k$
- Scattering: $V_k^r = V_k^C - V_k^i$
- Connection: $V_{k+1}^i = V_k^r$

Example (Cont'd)



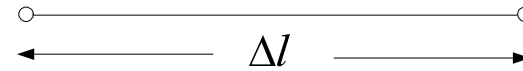
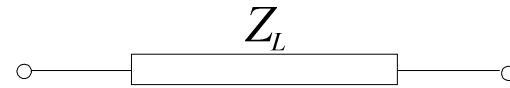
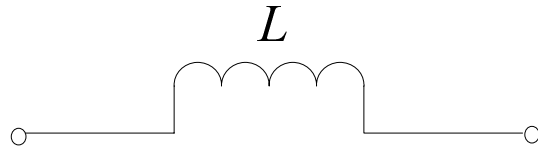
Example (Cont'd)



```
for k=1:10           %repeat for only time steps
    Ik=(Vs-2*Vi)/(R+Zc); %evaluate current
    CurrentValuesTLM(k,1)=Ik; %store current value
    CurrentValuesAnalytic(k,1)=(Vs/R)*exp(-(k-1)*dt/(R*C));
    Vk=2*Vi+Ik*Zc; %get capacitor voltage
    Vr=Vk-Vi; %evaluate reflected voltage (scattering)
    Vi=Vr; %evaluate new incident voltage (connection)
end
```

Try It!

Link Model of an Inductor



- Using similar derivation to that of the capacitor we have

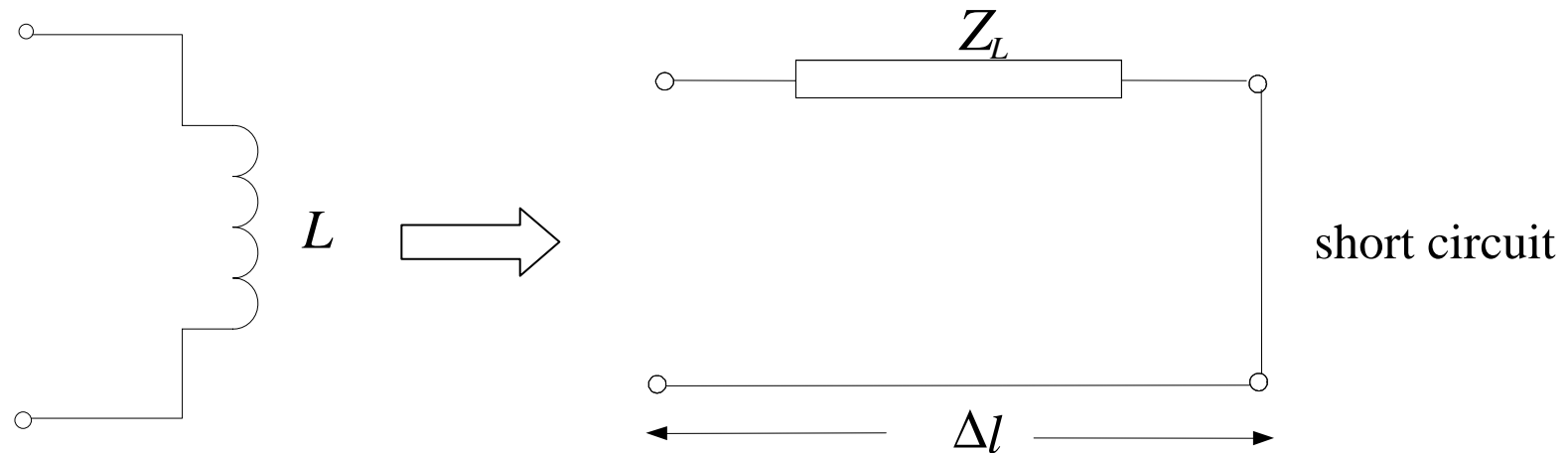
$$L_d \Delta l = L$$

$$u = \Delta l / \Delta t = 1 / \sqrt{L_d C_d} \implies C_d = \left(\frac{\Delta t}{\Delta l} \right)^2 \frac{1}{L_d}$$

$$Z_L = \sqrt{\frac{L_d}{C_d}} = \frac{L}{\Delta t}$$

$$\text{Error capacitance } C_e = C_d \Delta l = (\Delta t)^2 / L$$

Stub Model of an Inductor



- Using similar derivation we have

$$L_d \Delta l = L$$

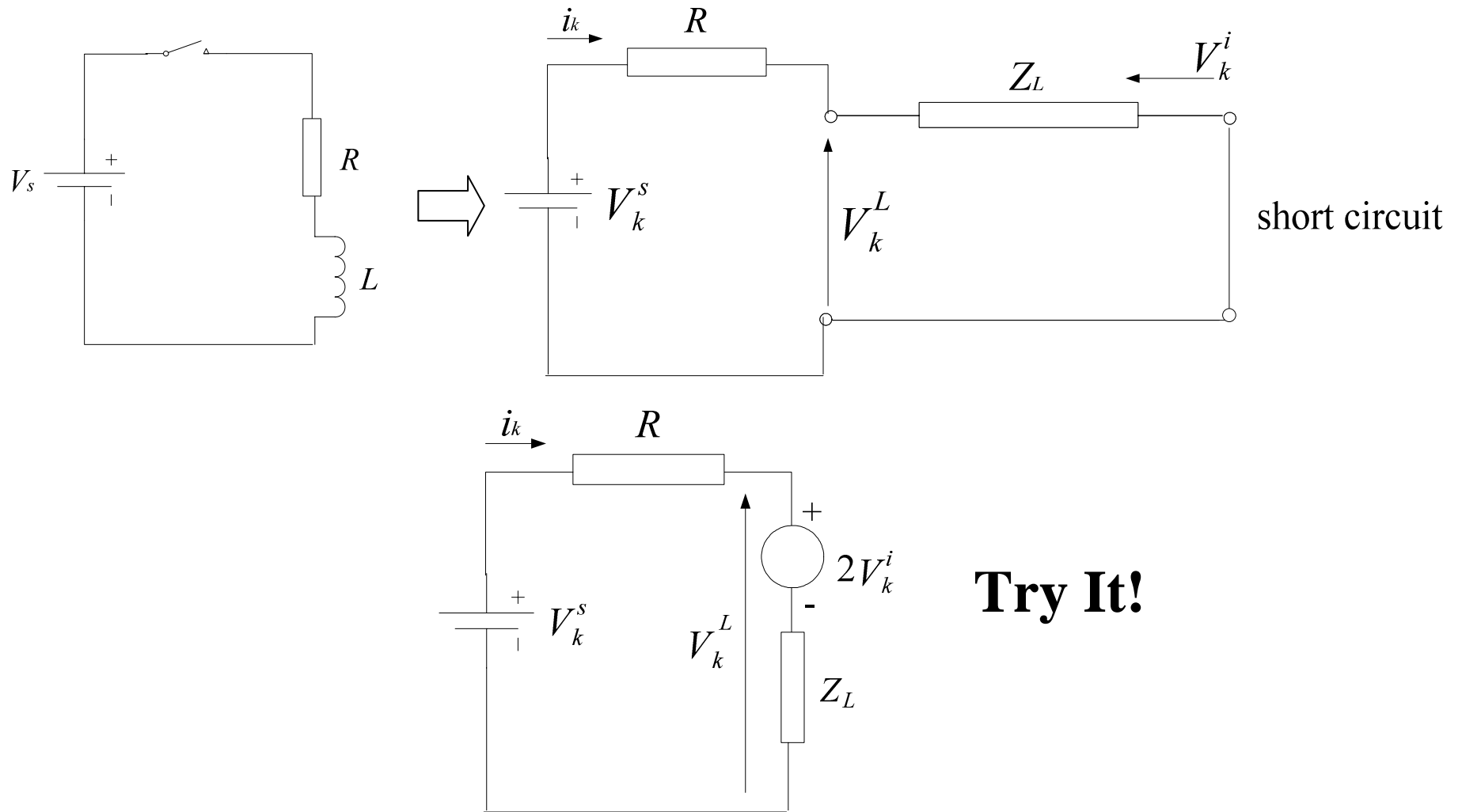
$$u = \Delta l / (\Delta t / 2) = 1 / \sqrt{L_d C_d} \implies C_d = \left(\frac{\Delta t}{\Delta l} \right)^2 \frac{1}{4L_d}$$

$$Z_L = \sqrt{\frac{L_d}{C_d}} = \frac{2L}{\Delta t}$$

$$\text{Error capacitance } C_e = C_d \Delta l = (\Delta t)^2 / (4L)$$

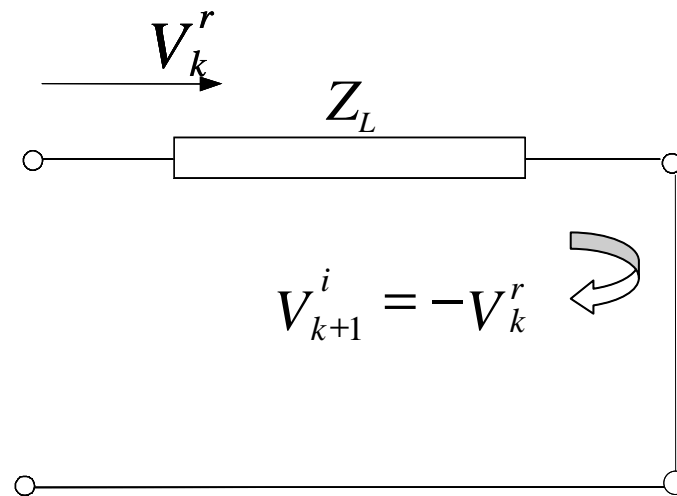
Example

- Evaluate the transient response of the shown circuit



Try It!

Example (Cont'd)



- Evaluate the current $i_k = (V_k^S - 2V_k^i) / (Z_L + R)$
- Evaluate the inductor voltage $V_k^C = 2V_k^i + Z_L i_k$
- Scattering: $V_k^r = V_k^C - V_k^i$
- Connection: $V_{k+1}^i = -V_k^r$