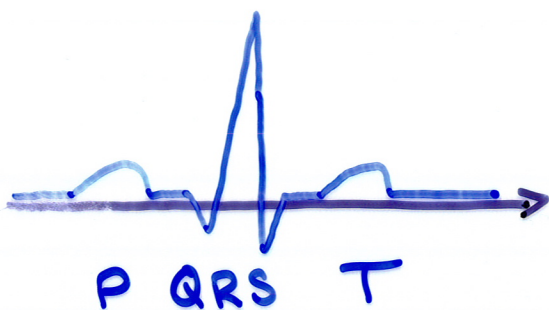
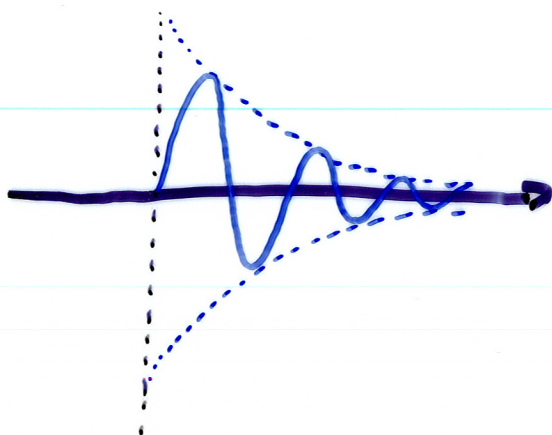


## A MOTIVATING EXAMPLE



To build a defibrillator, we would like to mimic this waveform

Let's try:  $e^{-\sigma t} \sin \omega t$

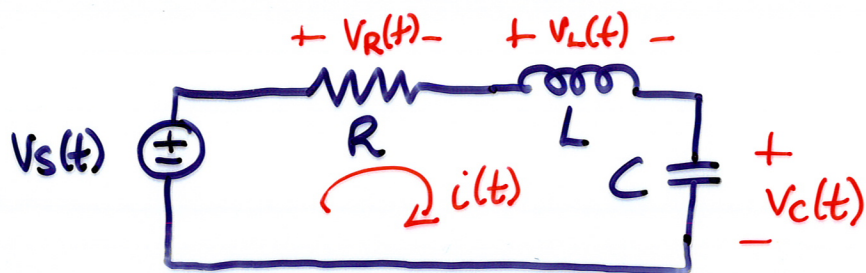


Not too bad.

Furthermore, we know how to synthesize this waveform,  
in a portable device!

Well, after this lecture we will know

## MORE FORMAL ANALYSIS



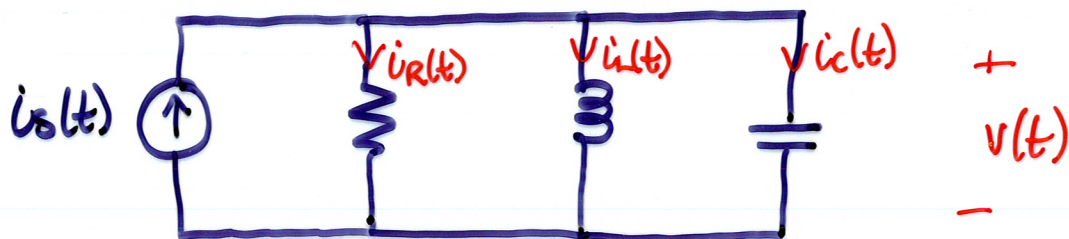
$$-V_s(t) + V_R(t) + V_L(t) + V_C(t) = 0$$

$$\Rightarrow Ri(t) + L \frac{di(t)}{dt} + \frac{1}{C} \int_{-\infty}^t i(\tau) d\tau = V_s(t)$$

Differentiate

$$R \frac{di(t)}{dt} + L \frac{d^2i(t)}{dt^2} + \frac{1}{C} i(t) = \frac{dV_s(t)}{dt}$$

$$\Rightarrow \frac{d^2i(t)}{dt^2} + \frac{R}{L} \frac{di(t)}{dt} + \frac{1}{LC} i(t) = \frac{1}{L} \frac{dV_s(t)}{dt}$$



KCh  $i_s(t) = i_R(t) + i_L(t) + i_C(t)$

$$i_s(t) = \frac{v(t)}{R} + \frac{1}{L} \int_{-\infty}^t v(x) dx + C \frac{dv(t)}{dt}$$

Differentiate and divide by C

$$\frac{d^2 v(t)}{dt^2} + \frac{1}{RC} \frac{dv(t)}{dt} + \frac{1}{LC} v(t) = \frac{1}{C} \frac{di_s(t)}{dt}$$

## SOLUTION OF SECOND ORDER DIFFERENTIAL EQUATIONS

$$\bullet \frac{d^2 x(t)}{dt^2} + a_1 \frac{dx(t)}{dt} + a_0 x(t) = f(t)$$

As in the first order case, natural and forced components

$$x(t) = x_n(t) + x_f(t)$$

### NATURAL RESPONSE

Solution when  $f(t) = 0$ ; ie.

$$\bullet \frac{d^2 x(t)}{dt^2} + a_1 \frac{dx(t)}{dt} + a_0 x(t) = 0 \quad \text{(*)}$$

⇒ need a function that has the same shape as its first and second derivatives

The exponential  $Ae^{st}$  looks like a good candidate

Here we will ~~consider~~ allow  $s$  to be real-valued or complex-valued.

Let's try. Substitute  $Ae^{st}$  into (\*)

$$As^2e^{st} + a_1 Ase^{st} + a_0 Ae^{st} = 0.$$

$$\Rightarrow Ae^{st} (s^2 + a_1 s + a_0) = 0$$

- $A=0$  is a solution, but is unlikely to satisfy initial conditions
- Therefore we seek solutions to the "characteristic equation"

$$s^2 + a_1 s + a_0 = 0$$

- Using standard formula, solutions are

$$s_1, s_2 = \frac{-a_1 \pm \sqrt{a_1^2 - 4a_0}}{2}$$

- Three regimes

$$* a_1^2 - 4a_0 > 0 : \text{overdamped}$$

$$* a_1^2 - 4a_0 = 0 : \text{critically damped}$$

$$* a_1^2 - 4a_0 < 0 : \text{underdamped.}$$

## SUPERPOSITION

- In general there are two solutions to the characteristic equation
- Which one should we use?
- Since the differential equation is linear, we can use both

## SIMPLIFYING CONCEPTS

Equation:  $s^2 + a_1 s + a_0 = 0$

Define

undamped natural frequency:  $\omega_0 = \sqrt{a_0}$

damping ratio  $\zeta = \frac{a_1}{2\omega_0}$

Therefore:  $s_1, s_2 = -\zeta\omega_0 \pm \omega_0\sqrt{\zeta^2 - 1}$

and

$\zeta > 1 \Rightarrow$  overdamped.  $[s_1, s_2 \text{ real}]$

$\zeta = 1 \Rightarrow$  critically damped  $[s_1 = s_2]$

$\zeta < 1 \Rightarrow$  under damped  $[s_1, s_2 \text{ complex}]$

Overdamped:  $\zeta > 1$ ,  $s_1, s_2$  real

$$x_n(t) = k_1 e^{s_1 t} + k_2 e^{s_2 t}$$

$$\begin{aligned} s_1, s_2 &= -\zeta \omega_0 \pm \omega_0 \sqrt{\zeta^2 - 1} \\ &= \frac{-a_1 \pm \sqrt{a_1^2 - 4a_0}}{2} \end{aligned}$$

Underdamped:  $\zeta < 1$ ,  $s_1, s_2$  complex

Complex exponentials correspond to exponentially decaying sinusoids

$$x_n(t) = A_1 e^{-\sigma t} \cos(\omega_d t) + A_2 e^{-\sigma t} \sin(\omega_d t)$$

where

$$\sigma = \zeta \omega_0$$

$$\omega_d = \omega_0 \sqrt{1 - \zeta^2} \quad \text{damped natural frequency}$$

Critically damped:  $\zeta = 1$   $s_1 = s_2 = -\zeta \omega_0 = \sigma$

$$x_n(t) = B_1 e^{-\sigma t} + B_2 t e^{-\sigma t}$$

Finding constants: once the form of  $x_f(t)$  also known

Constants typically found from

$$x(t)|_{t=0};$$

$$\frac{dx(t)}{dt}|_{t=0}; \text{ eqn itself};$$

if  $f(t)$  is constant also  $x(\infty)$