

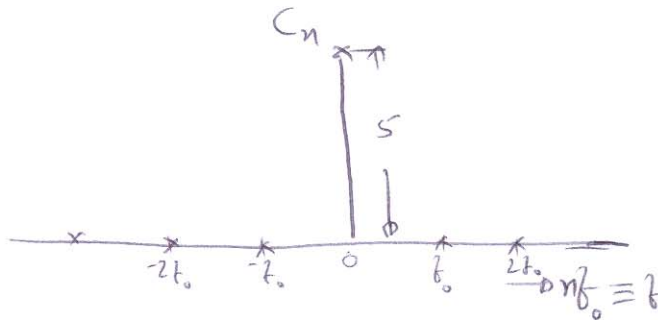
FOURIER TRANSFORM OF PERIODIC SIGNALS

EXAMPLE 1:

CONSIDER A DC SIGNAL

$$g_p(t) = 5 \text{ V}$$

FOURIER COEFFICIENTS: $C_0 = 5$, ~~where $C_n = 0$~~
 $C_n = 0$ IF $n \neq 0$



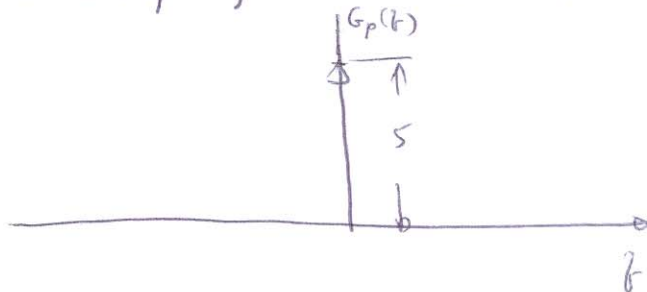
FOURIER TRANSFORM:

$$\delta(t) \rightleftharpoons 1$$

DUALITY PROPERTY $\rightarrow 1 \rightleftharpoons \delta(f)$

$$5 \rightleftharpoons 5\delta(f)$$

$$\therefore F[g_p(t)] = 5\delta(f) = G_p(f)$$



COMPARING FIGS. 1 & 2, WE SEE THAT, $G_p(f) = C_0 \delta(f)$.

(2)

EXAMPLE 2 :

$$g_p(t) = 5 \cos(2\pi 500 t)$$

FOURIER SERIES : $C_0 = 0$ (\because NO DC COMPONENT)

$$f_0 = 500, \quad C_1 = \frac{5}{2} \quad (\because \text{WEIGHT OF } \exp(j 2\pi 500 t))$$

$$C_{-1} = \frac{5}{2} \quad (\because \text{WEIGHT OF } \exp(-j 2\pi 500 t))$$

$$C_n = 0 \quad \text{IF } n \neq 1.$$

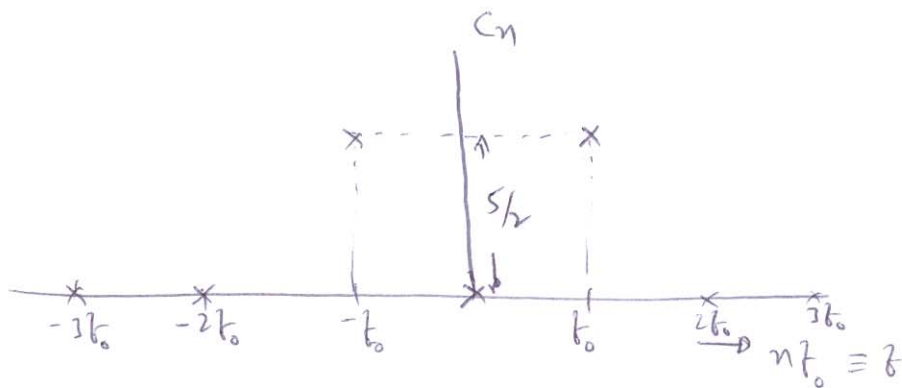


FIG. 3

FOURIER TRANSFORM :

$$\text{expt. } 1 \equiv \delta(\tau)$$

$$\exp(j 500 \cdot 2\pi t) \equiv \delta(\tau - 500) \rightarrow (1)$$

$$\exp(-j 2\pi \cdot 500 t) \equiv \delta(\tau + 500) \rightarrow (2)$$

$$\text{ADDING (1) \& (2), } \cos(2\pi 500 t) \equiv \frac{1}{2} \delta(\tau - 500) + \frac{1}{2} \delta(\tau + 500)$$

$$5 \cos(2\pi 500 t) \equiv \frac{5}{2} \delta(\tau - 500) + \frac{5}{2} \delta(\tau + 500)$$

$$(\because \cos(\theta) = \frac{\exp(j\theta) + \exp(-j\theta)}{2}, \theta = 2\pi 500 t)$$

(3)

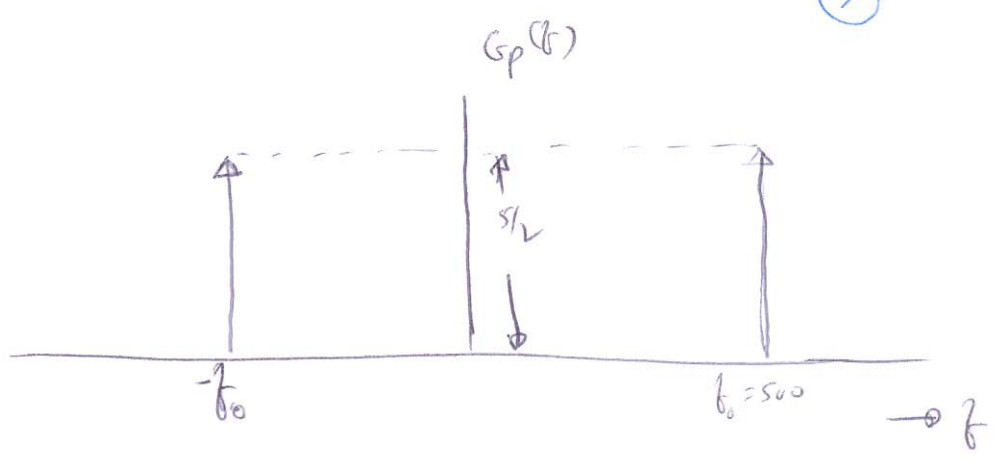


FIG. 4.

FMM FIGS. 3 & 4, WE SEE THAT

$$G_p(f) = C_1 \delta(f - f_0) + C_2 \delta(f + f_0)$$

EXAMPLE 3:

$$g_p(t) = 5 \cos(2\pi 500t) + 2 \cos(2\pi 1000t)$$

FOURIER SERIES:

$f_0 = 500 \text{ Hz}$
 $2f_0 = 1000 \text{ Hz}$

- $C_1 = 5/2$ (FUNDAMENTAL)
- $C_1 = 5/2$ (")
- $C_2 = 2/2$ (SECOND HARMONIC)
- $C_2 = 2/2$ (")
- $C_n = 0$ IF $n \neq 1$ OR $n \neq 2$

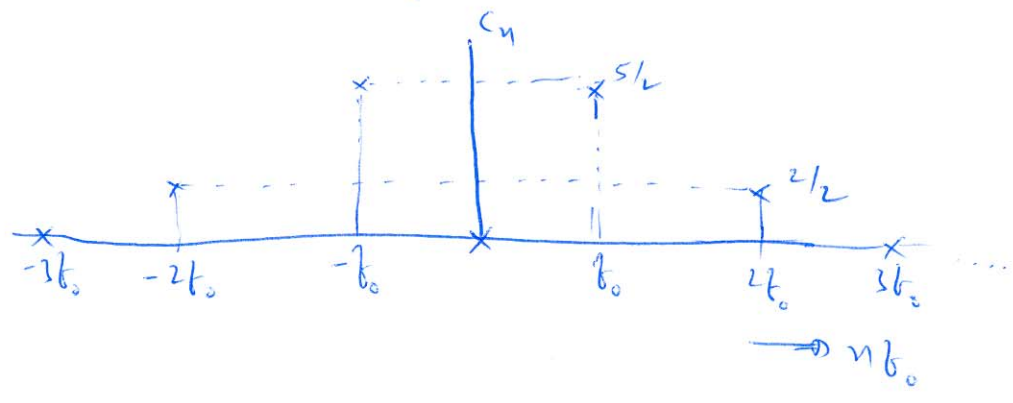


FIG. 5

(4)

FOURIER SPECTRUM IS A PLOT OF THE ^(EOT) AMPLITUDES OF
(ACTUALLY THAT OF $\exp(j2\pi ft)$)
SINE OR COSINE WAVES AS A FUNCTION OF THE FREQUENCY.

FOURIER TRANSFORM:

$$5 \cos(2\pi 500t) \implies \frac{5}{2} \delta(f-500) + \frac{5}{2} \delta(f+500) \rightarrow (3)$$

$$2 \cos(2\pi 1000t) \implies \frac{2}{2} \delta(f-1000) + \frac{2}{2} \delta(f+1000) \rightarrow (4)$$

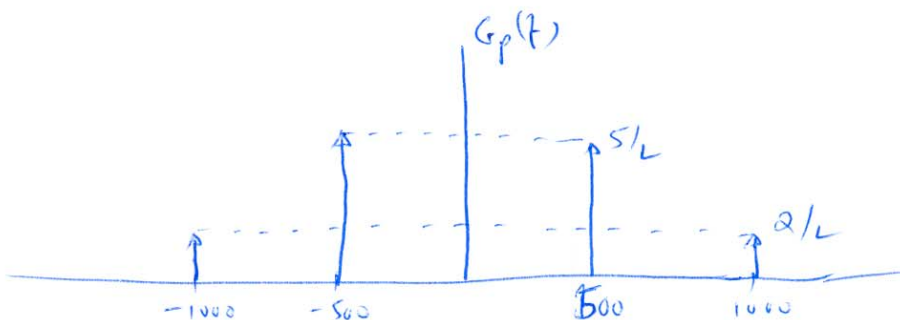


FIG. 6

$$G_p(f) = \frac{5}{2} \delta(f-500) + \frac{5}{2} \delta(f+500) + \frac{2}{2} \delta(f-1000) + \frac{2}{2} \delta(f+1000) \rightarrow (5)$$

$$= \sum_{n=-\infty}^{\infty} c_n \delta(f - n f_0)$$

$$= c_1 \delta(f-500) + c_{-1} \delta(f+500) + c_2 \delta(f-1000) + c_{-2} \delta(f+1000) \rightarrow (6)$$

NOTE THE SIMILARITY BETWEEN FOURIER SERIES SPECTRUM (FIG. 5)
& FOURIER TRANSFORM SPECTRUM (FIG. 6).

(5)

Eq. (6) can be written as

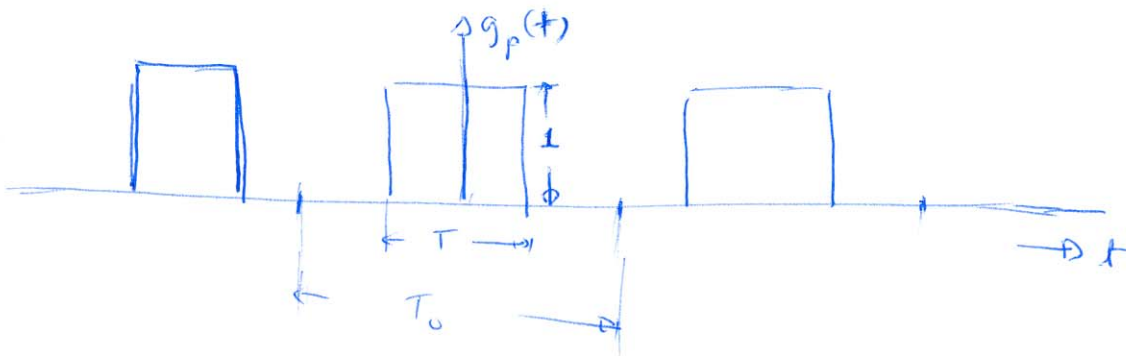
$$G_p(\omega) = \sum_{n=-2}^2 C_n \delta(\omega - n\omega_0), \quad \omega_0 = 500 \text{ Hz}$$

In general, if a periodic signal $g_p(t)$ has a Fourier coefficients C_n , its Fourier transform can be written as

$$G_p(\omega) = \sum_{n=-\infty}^{\infty} C_n \delta(\omega - n\omega_0)$$

EXAMPLE 4:

CONSIDER A 50% DUTY CYCLE SQUARE WAVE



$$\text{DUTY CYCLE} = \frac{\text{PULSE WIDTH}}{\text{PERIOD}} \times 100 = \frac{T}{T_0} \times 100$$

$$\therefore 50\% \text{ DUTY CYCLE} \Rightarrow T = T_0/2$$

(6)

$$\begin{aligned}
 C_n &= \frac{1}{T_0} \int_{-T_0/2}^{T_0/2} g_p(t) \exp(-j2\pi n f_0 t) dt, \quad f_0 = 1/T_0 \\
 &= \frac{1}{T_0} \int_{-T/2}^{T/2} 1 \cdot \exp(-j2\pi n f_0 t) dt \\
 &= \frac{1}{T_0} \left. \frac{e^{-j2\pi n f_0 t}}{-j2\pi n f_0} \right|_{-T/2}^{T/2} = \frac{1}{2} \text{SINC}(n/2) \\
 &\quad \text{(USE } T = T_0/2)
 \end{aligned}$$

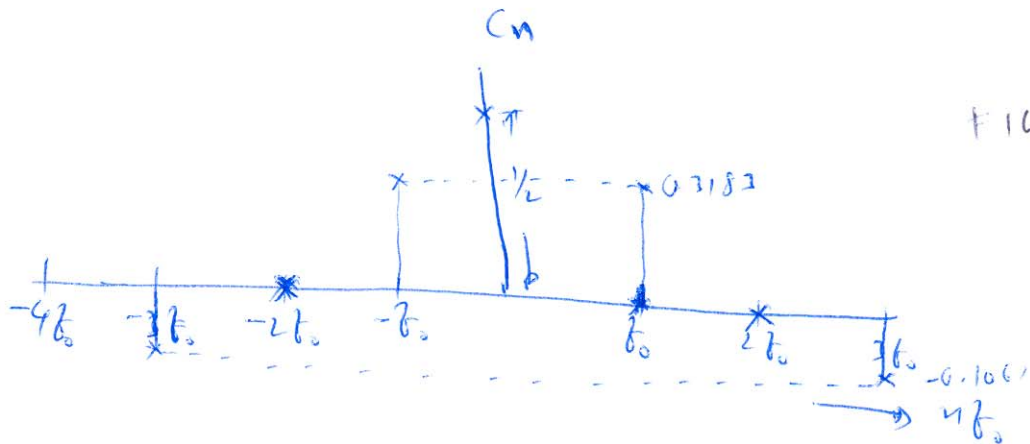
$$\text{SINC}(x) = \frac{\sin(\pi x)}{\pi x}$$

DC component, $C_0 = \frac{1}{2} \sin(0) = \frac{1}{2}$

$$C_1 = \frac{1}{2} \text{SINC}(1/2) = 0.3183$$

$$C_2 = \frac{1}{2} \text{SINC}(2/2) = 0 \quad (\because \text{SINC}(n) = 0, \quad n = \text{integer} \neq 0)$$

$$C_3 = \frac{1}{2} \text{SINC}(3/2) = -0.1061$$



FOURIER TRANSFORMS:

$$F[g_p(t)] = G_p(f) = \sum_{n=-\infty}^{\infty} c_n \delta(f - n f_0)$$

$$= \frac{1}{2} \sum_{n=-\infty}^{\infty} \text{sinc}(n/L) \delta(f - n f_0)$$

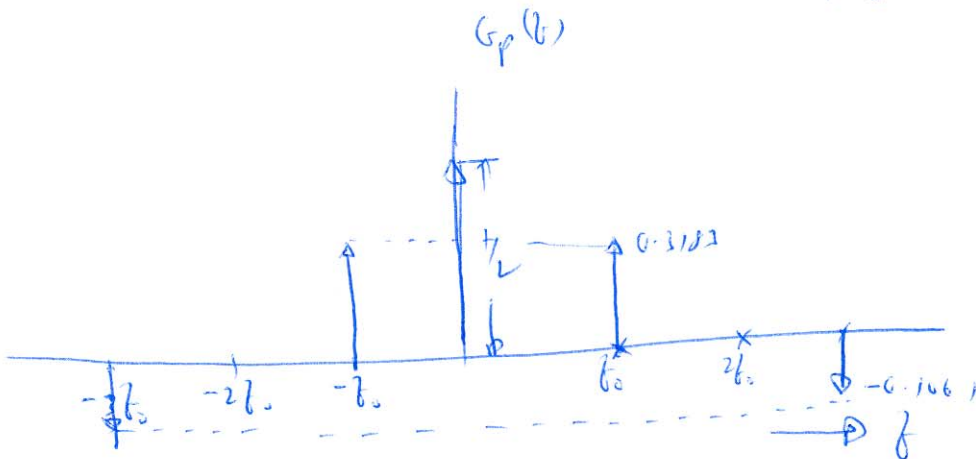
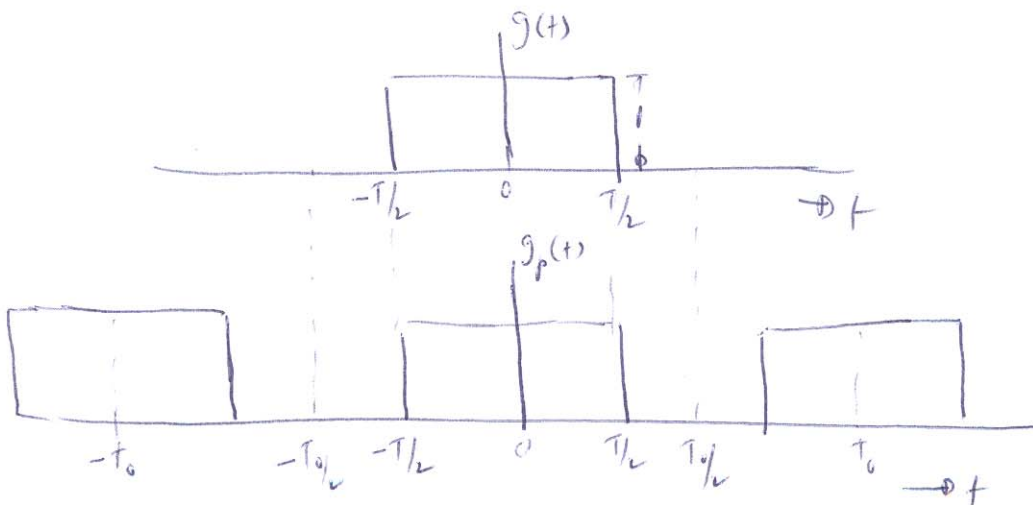


FIG. 8

EXAMPLES: IN EXAMPLE 4, CONSIDER THE RECTANGULAR PULSE WITHIN A PERIOD T_0 .

$$g(t) = \text{rect}(t/T)$$



$$g_p(t) = \sum_{n=-\infty}^{\infty} g(t - nT_0)$$

$$g_p(t) = \sum_{n=-\infty}^{\infty} \text{rect}\left(\frac{t - nT_0}{T}\right)$$

$g_p(t)$ IS PERIODIC & IT IS KNOWN AS 'PERIODIC EXTENSION OF $g(t)$ '.

(8)

$$\text{SINCE } \text{sinc}(t) \rightleftharpoons \text{SINC}(f)$$

$$\text{SCALING PROPERTY} \quad \text{sinc}(t/T) \rightleftharpoons T \text{SINC}(fT)$$

$$T = T_0/2$$

$$\therefore G(f) = T \text{SINC}(fT) = \frac{T_0}{2} \text{SINC}(fT_0/2)$$

CONSIDER THE FOURIER SPECTRUM OF $g(t)$ at n^{th} HARMONIC OF ~~T~~ ITS PERIODIC EXTENSION $g_p(t)$.

$$G(f) \Big|_{nf_0} = \frac{T_0}{2} \text{SINC}(nf_0 T_0/2) = \frac{T_0}{2} \text{SINC}(n/2) \quad (\because f_0 T_0 = 1)$$

SINCE THE FOURIER COEFFICIENTS OF $G_p(f)$ ARE

$$C_n = \frac{1}{2} \text{SINC}(n/2)$$

IT FOLLOWS THAT

$$G(f) \Big|_{nf_0} = G(nf_0) = \cancel{T_0} T_0 C_n \quad \rightarrow (*)$$

OR

$$\boxed{C_n = \frac{1}{T_0} G(nf_0)}$$

IN EXAMPLE 4, WE HAVE SEEN THAT

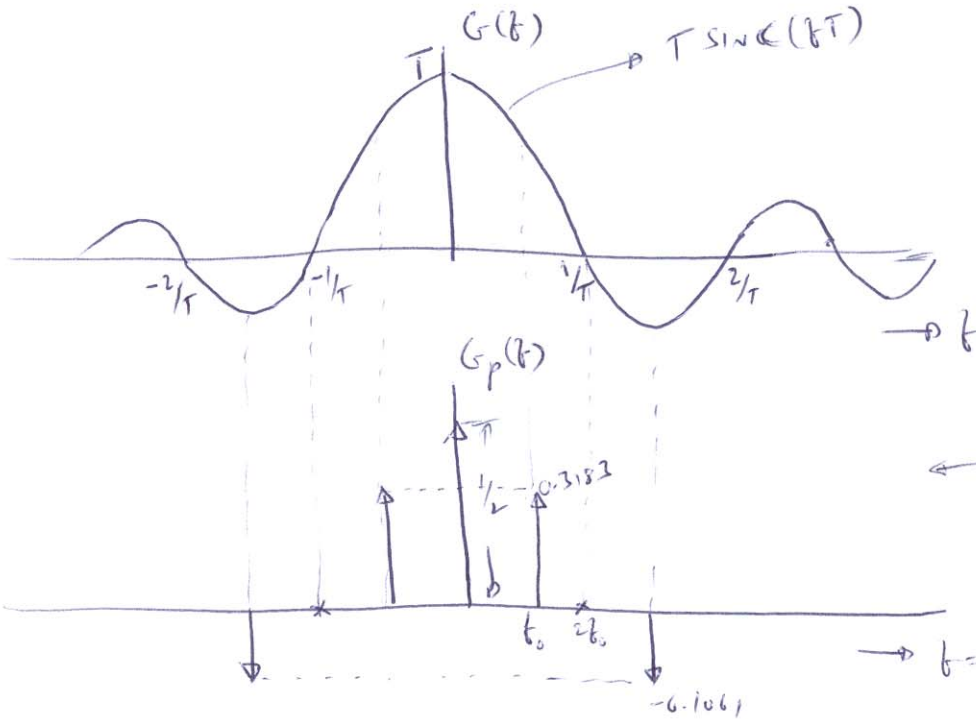
$$G_p(f) = \sum_{n=-\infty}^{\infty} C_n \delta(f - nf_0) \quad \rightarrow (**)$$

(9)

USING EQ. (*) IN (**), WE FIND

$$G_p(f) = T_0 \sum_{n=-\infty}^{\infty} G(nf_0) \delta(f - nf_0) \rightarrow (***)$$

ALTHOUGH WE PROVED THE ABOVE PROPERTY FOR A PERIODIC WAVE
 CONSISTING OF RECTANGULAR PULSES, IT HOLDS TRUE FOR ARBITRARY
 TIME LIMITED SIGNAL $g(t)$ WITHIN A PERIOD T_0 .



$$T = T_0/2$$

FIG. 9

← SAME AS FIG. 8

FIG. 10

MEANING OF EQ. (***)

SUPPOSE WE KNOW THE FOURIER TRANSFORM OF A PULSE $g(t)$ ($= G(f)$). ITS
 PERIODIC EXTENSION $g_p(t)$ WITH PERIOD T_0 CAN BE OBTAINED BY
 TAKING THE SAMPLES OF $G(f)$ AT $f = nT_0$ & MULTIPLYING THEM BY
 DELTA FUNCTIONS LOCATED AT $f = nT_0$, AS SHOWN IN FIG. 10.