

Lecture 3

From sections 5.8, 5.9 and 6.2

Solve 5.24, 5.26, 5.27, 5.30, 5.32,
5.33, 5.34

Self Read

Review Sections 5.5 - 5.7

* For air $\bar{D} = \epsilon_0 \bar{E}$

* For dielectric $\bar{D} = \epsilon_0 \bar{E} + \bar{P}$

* For a linear dielectric $\bar{D} = \chi_e \epsilon_0 \bar{E}$

\Downarrow
 $\bar{D} = \epsilon_0 \Sigma_r \bar{E}$ where

$$(e_x + 1) = r^3$$



Continuity Equation

I_{out}

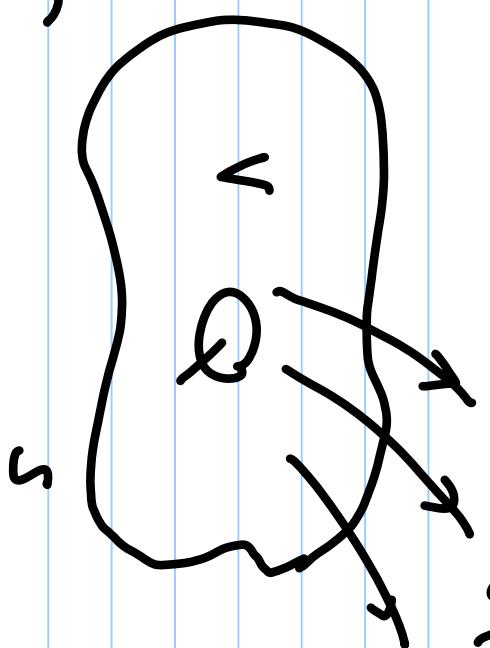
$$* \int_{out} = \oint \vec{J} \cdot d\vec{s}$$

* Flow of current is

the outward direction

means that Q is decreasing

$$\oint \vec{J} \cdot d\vec{s} = - \frac{dQ}{dt} \quad \boxed{\int \int g_v dv}$$



Continuity Equation (Cont'd)

* Using Divergence theorem

$$\oint \vec{J} \cdot d\vec{r} = \iiint (\nabla \cdot \vec{J}) dv = - \frac{\partial}{\partial t} \iiint \rho v dv$$

$$\downarrow \quad \nabla \cdot \vec{J} = - \frac{\partial \rho v}{\partial t}$$

* For steady currents, $-\frac{\partial \rho v}{\partial t} = 0$
 $\nabla \cdot \vec{J} = 0$ (Kirchhoff's Current Law)

Relaxation Time

* For a linear medium $\bar{J} = \sigma E$

$$\downarrow \nabla \cdot (\sigma E) = - \frac{\partial g_V}{\partial x}$$

$$\downarrow \frac{\partial}{\sum} (\nabla \cdot \epsilon E) = - \frac{\partial g_V}{\partial t}$$

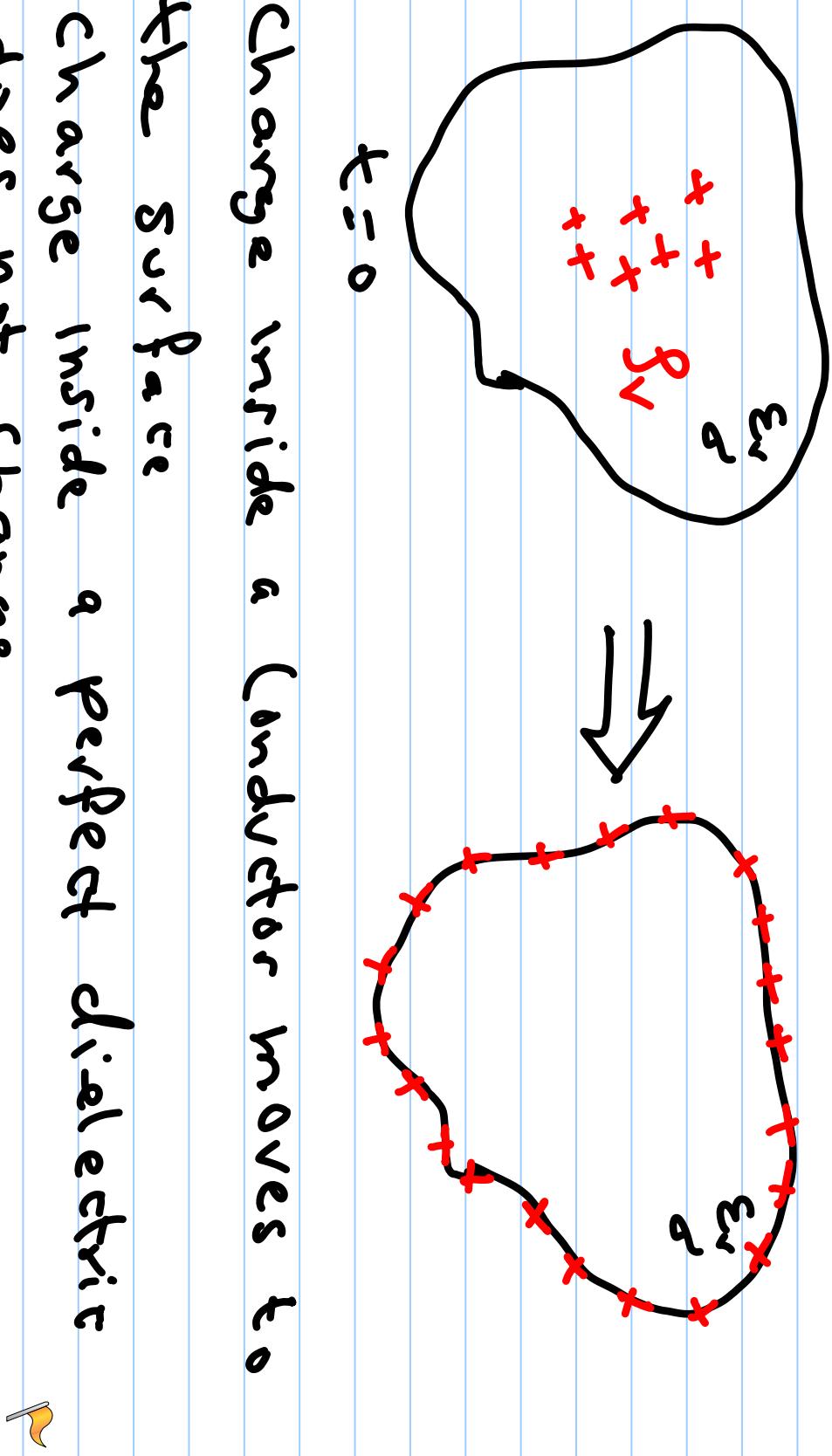
$$\frac{\partial g_V}{\partial x} + \frac{\partial}{\sum} g_V = 0 \quad (1st \text{ order D.E.})$$

$$\downarrow g_V = g_{V_0} e^{-t/T_R} \quad T_R = \frac{\sum}{\sigma} =$$

relaxation time

Relaxation Time (Cont'd)

- * charge inside a conductor moves to the surface
- * charge inside a perfect dielectric does not change



Boundary Conditions

* Dielectric - Dielectric Interface

$$\epsilon_1 E_1 = \epsilon_2 E_2$$

$$D_{1n} - D_{2n} = \rho_r$$

$$D_{1n} = D_{2n}$$

$$\epsilon_{r1} E_{1n} = \epsilon_{r2} E_{2n}$$



* Tangential component of \vec{E} is

continuous, while normal component of \vec{E} is

E_n is discontinuous

Perfect Conductor-Dielectric Interface



$$E_{1n} = E_{2n} = 0$$

$$D_{1n} - D_{2n} = \rho_c$$

$$D_{1n} = \rho_c$$

$$\rho_c \cdot \epsilon_2 = 0$$

- * only a normal E field component exists at the surface.

Laplace & Poisson's Equations

$$\nabla \cdot \vec{D} = \rho_v \rightarrow \nabla \cdot (\epsilon \vec{E}) = \rho_v$$

$$\nabla \cdot (\epsilon \nabla V) = -\rho_v$$

If ϵ is a constant, we have

$$\epsilon (\nabla \cdot \nabla V) = -\rho_v$$

$$\nabla^2 V = -\frac{\rho_v}{\epsilon} \leftarrow \text{Laplace equation}$$

$$\text{if } \rho_v = 0, \nabla^2 V = 0$$

Laplace & Poisson Equations

- * The Laplacian operator is given

$$\Delta^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}$$

$$\Delta^2 V = \frac{1}{g} \frac{\partial}{\partial g} \left(g \frac{\partial V}{\partial g} \right) + \frac{1}{g^2} \frac{\partial^2 V}{\partial \varphi^2} + \frac{\partial^2 V}{\partial z^2}$$

$$\Delta^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \theta^2} + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \varphi^2}$$

+ $\frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \varphi^2}$

