

Dr. Mohamed Bakr, EE3FK4, 2008

Note Title

2/12/2008

Lecture 10

From Sections 10.1 - 10.2

Solve 10.1

Wave Equations

- * We assume first that the electric field has only one component and it is a function only of $\Sigma \Rightarrow E = E_x g_x$
- * Using Maxwell's equation

$$\nabla \times \underline{E} = - \frac{\partial \underline{B}}{\partial t} = - \mu \epsilon \frac{\partial \underline{H}}{\partial t}$$
$$\rightarrow -\mu \epsilon \frac{\partial H_y}{\partial t} = \frac{\partial E_x}{\partial z} g_y \rightarrow H_y = H_0 g_y$$

Wave Equation (Cont'd)

* Also $\nabla \times H = J + \frac{\partial D}{\partial t} \rightarrow$ For a lossless medium with no sources we have

$$\nabla \times H = -\frac{\partial D}{\partial t} \rightarrow -\frac{\partial H_y}{\partial z} \text{ at } x = \epsilon \frac{\partial E_x}{\partial t} = \epsilon \frac{\partial E_x}{\partial t} \text{ at } x$$

$$\frac{\partial E_x}{\partial z} = -\mu \frac{\partial H_y}{\partial x} - \textcircled{1} \quad \epsilon \frac{\partial E_x}{\partial t} = -\frac{\partial H_y}{\partial z} - \textcircled{2}$$

2 First order PDE's in two variables that can be reduced to one second order equation in one variable!

Wave Equation (Cont'd)

Differentiate (1) relative to z and (2)
relative to t and eliminate $\frac{\partial^2 H_0}{\partial z \partial t}$ to

get $\frac{\partial^2 E_x}{\partial z^2} = \mu \frac{\partial^2 E_x}{\partial t^2}$

Define $v_p = \sqrt{\mu \epsilon}$ as the phase velocity

Wave eqn: $\frac{\partial^2 E_x}{\partial z^2} - \frac{1}{v_p^2} \frac{\partial^2 E_x}{\partial t^2} = 0$

Solution of Wave equation

The solution of this equation is

$$E_x = f_1\left(t - \frac{z}{v_p}\right) + f_2\left(t + \frac{z}{v_p}\right)$$

Wave travelling
in $+z$ direction

Wave travelling
in $-z$ direction

Example : $f_1\left(t - \frac{z}{v_p}\right) = \cos\left(t - \frac{z}{v_p}\right)$

$$f_1\left(t - \frac{z}{v_p}\right) = e^{-(t - \frac{z}{v_p})}$$

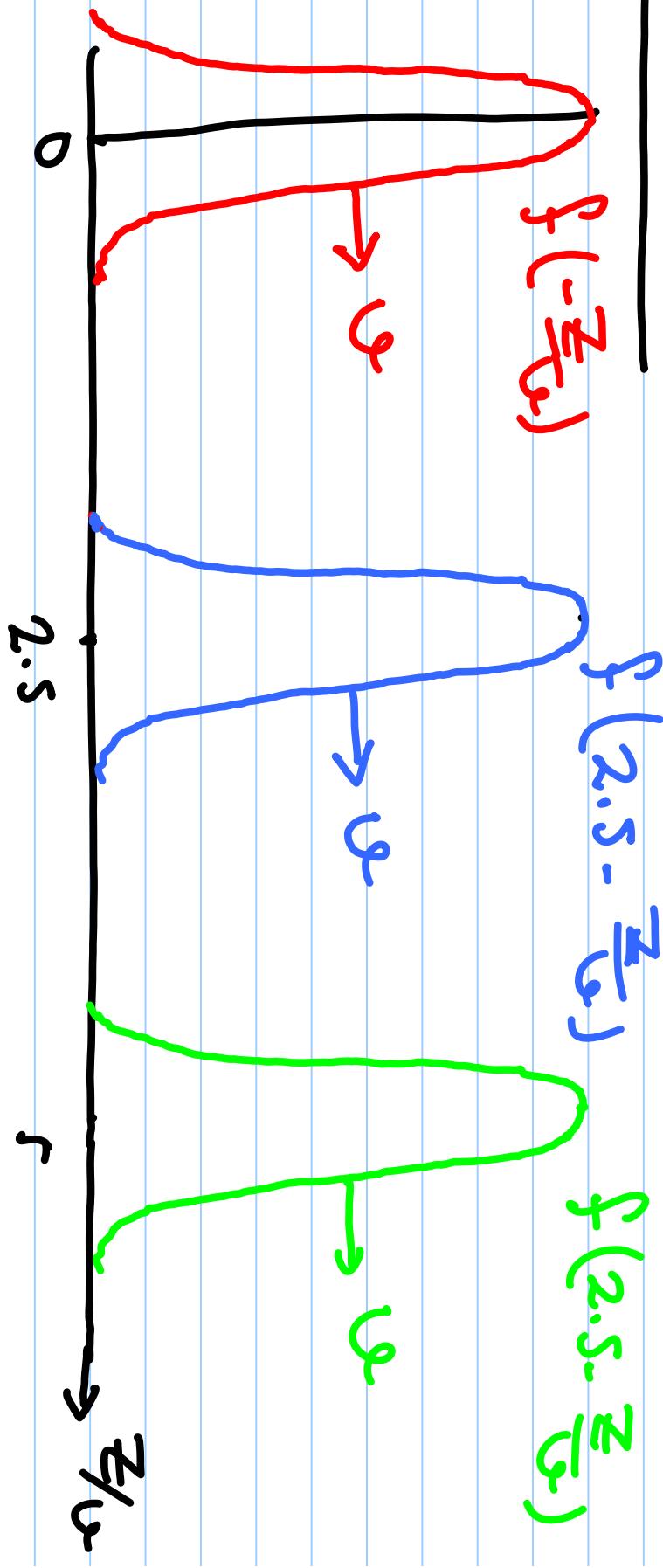
1D wave

$$f(-\frac{\pi}{\omega})$$

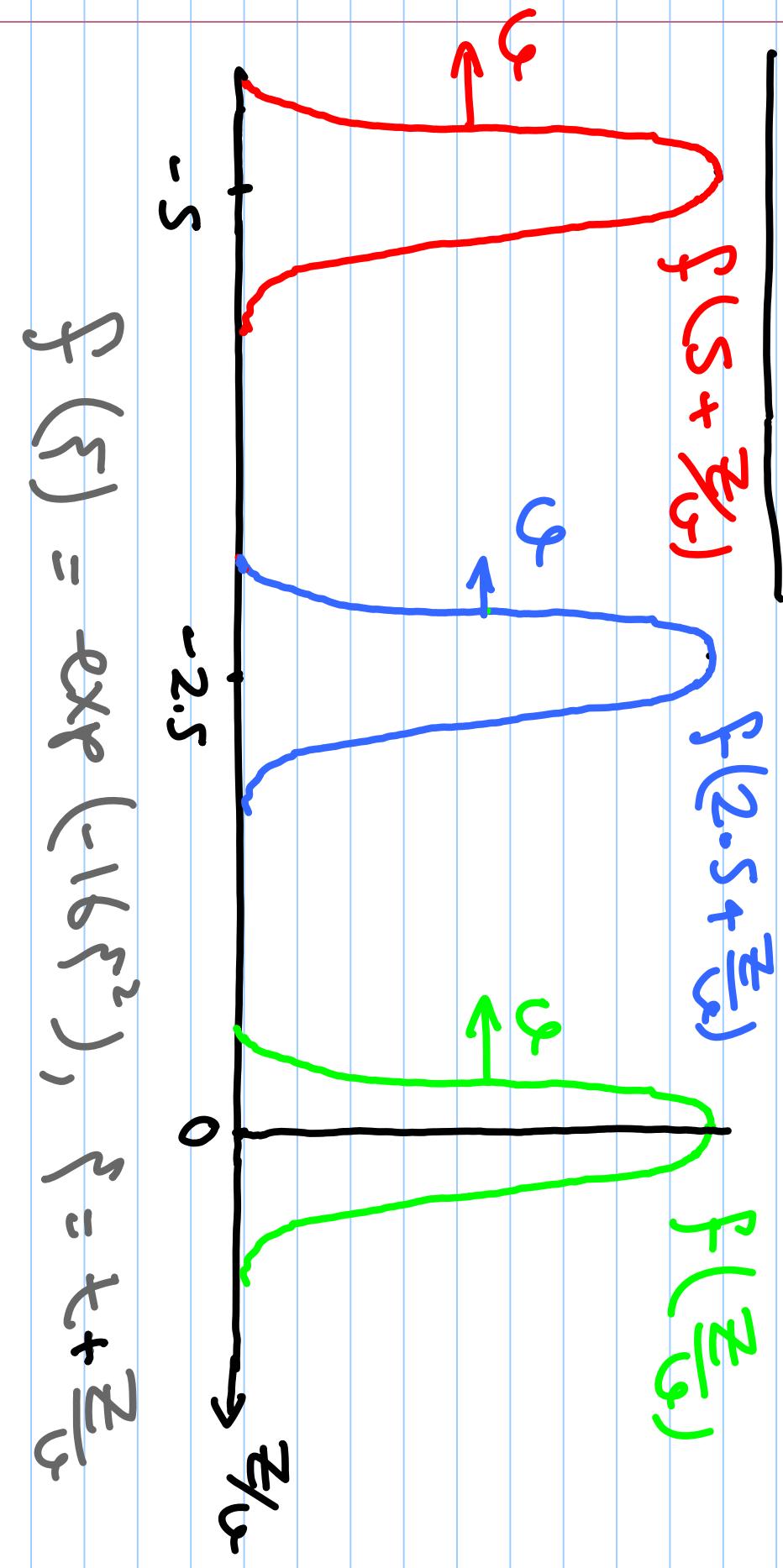
$$f(2.5 - \frac{\pi}{\omega})$$

$$f(2.5 - \frac{\pi}{\omega})$$

$$f(\xi) = \exp(-16\xi^2), \quad \xi = \tau - \frac{\pi}{\omega}$$



1D wave (cont'd)



$$f(\xi) = \exp(-16\xi^2), \quad \xi = \zeta + \frac{z}{\omega}$$

The Sineoidal Case

$$f(t) = \cos(\omega t + \phi)$$

$$\downarrow \\ f\left(t - \frac{\pi}{\omega}\right) = \cos\left(\omega\left(t - \frac{\pi}{\omega}\right) + \phi\right)$$

$$\downarrow \\ f\left(t - \frac{\pi}{\omega}\right) = \cos\left(\omega t - \frac{\pi}{\omega} + \phi\right)$$

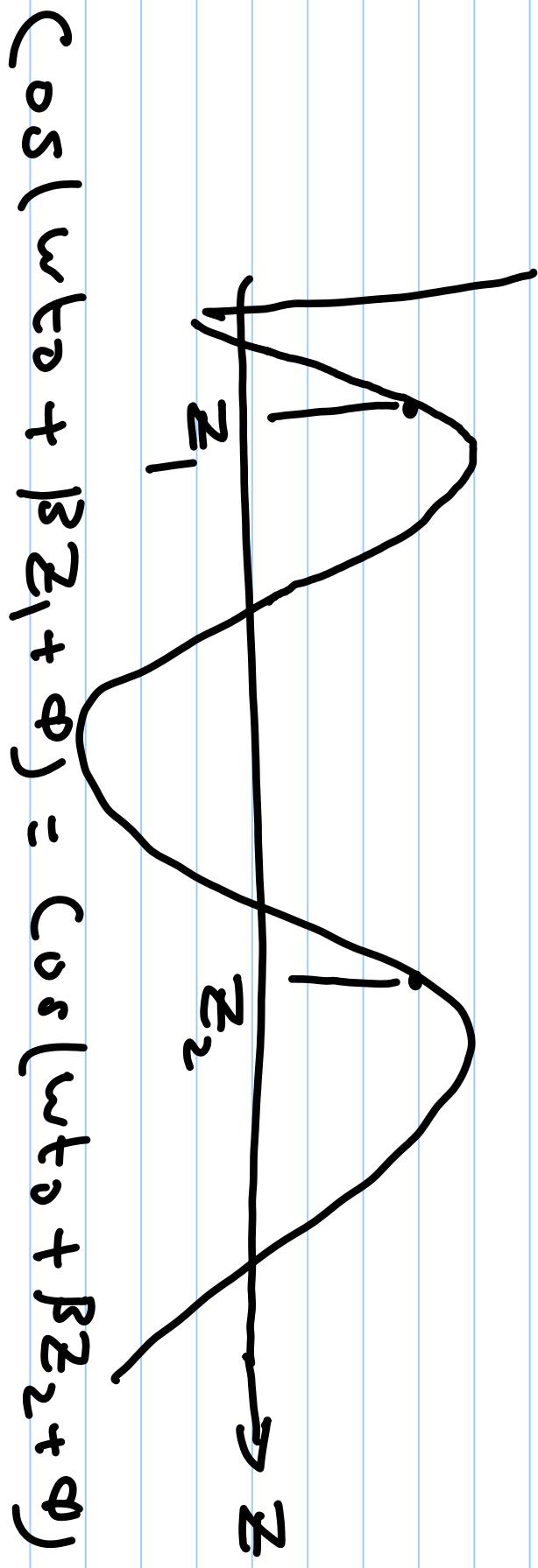
$$\downarrow \\ f\left(t - \frac{\pi}{\omega}\right) = \cos\left(\omega t - \pi + \phi\right)$$

$$\text{Similarly, } f\left(t + \frac{\pi}{\omega}\right) = \cos\left(\omega t + \pi + \phi\right)$$

The Sinusoidal Case (Cont'd)

$\beta = \text{Propagation Constant} = \frac{\omega}{c}$

* If we fix time, we have



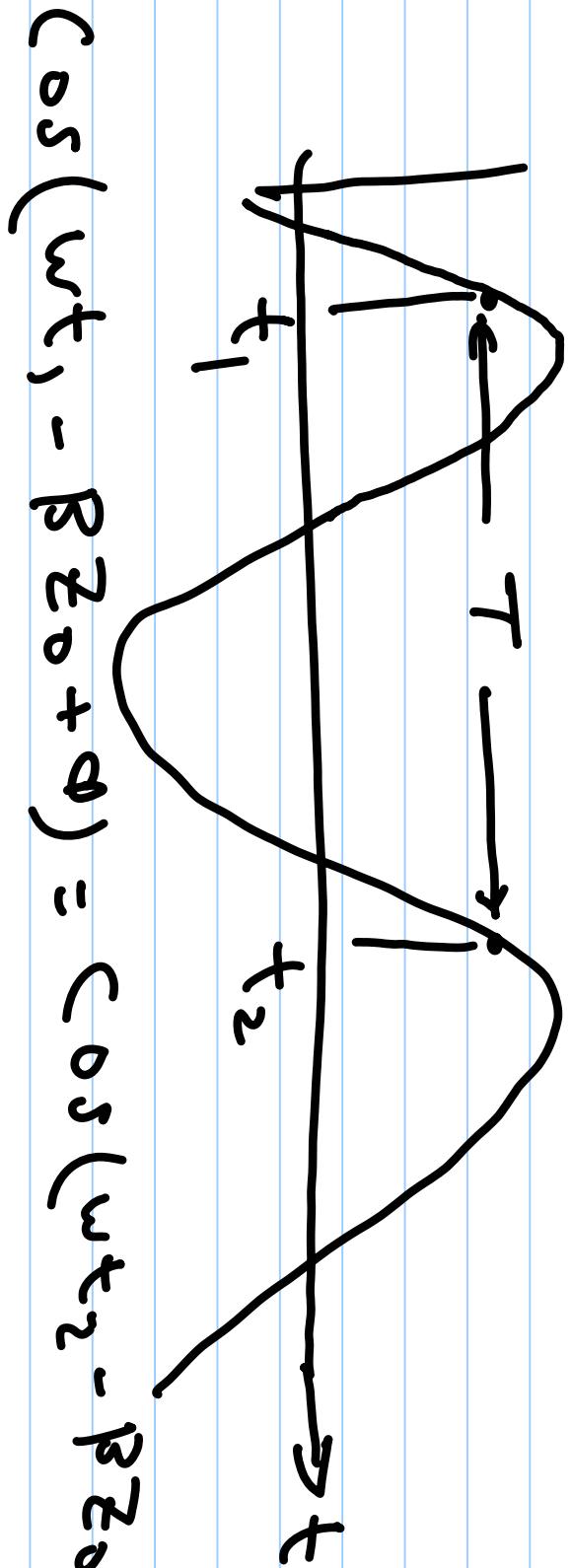
$$\cos(\omega t_0 + \beta \bar{z}_1 + \phi) = \cos(\omega t_0 + \beta \bar{z}_2 + \phi)$$

The Sinusoidal Case ((cont'd))

$$\Rightarrow \beta(z_r - z_1) = 2\pi$$

$$\beta\gamma = 2\pi, \quad \gamma = \frac{2\pi}{\beta} = \text{wavelength}$$

* If we fix z , we have



$$\cos(\omega t, -\beta z_0 + \phi) = \cos(\omega t_2 - \beta z_0 + \phi)$$

The Sineoidal Case (Cont'd)

$$\omega(t_2 - t_1) = 2\pi$$

$\overbrace{}^T$

$$\omega T = 2\pi \rightarrow T = \frac{2\pi}{\omega} = \frac{1}{f}$$

T = periodic time

