

Dr. Mohamed Bakr, EE3FK4, 2008

Note Title

2/7/2008

LECTURE 9

From Sections 9.5 - 9.7 of textbook

Solve 9.18, 9.21, 9.23, 9.27, 9.32, 9.37,

9.38, 9.39

Maxwell's Equations

Integral Form

$$\oint \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \iint \vec{B} \cdot d\vec{s}$$

$$\oint \vec{H} \cdot d\vec{l} = \iint (\vec{J} + \frac{\partial \vec{D}}{\partial t}) \cdot d\vec{s}$$

$$\oint \vec{D} \cdot d\vec{s} = \iiint \rho_v dv$$

$$\oint \vec{B} \cdot d\vec{s} = 0$$

$$\oint \vec{J} \cdot d\vec{s} = -\frac{\partial}{\partial t} \iiint \rho_v dv$$

Differential Form

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

$$\nabla \cdot \vec{D} = \rho_v$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \cdot \vec{J} = -\frac{\partial \rho_v}{\partial t}$$

Constitutive Equations

* The quantities \underline{E} , \underline{D} , \underline{H} , \underline{B} , and \underline{J} are related through material properties

$$* \underline{D} = \epsilon \underline{E} = \epsilon_0 \epsilon_r \underline{E}$$

$$\underline{B} = \mu \underline{H} = \mu_0 \mu_r \underline{H}$$

$$\underline{J} = \sigma \underline{E}$$

In general ($\underline{J}_t = \sigma \underline{E} + \underline{J}_i$)

Boundary Conditions

* Same boundary conditions still apply at interfaces of different materials.

$$E_{1t} - E_{2t} = 0 \Rightarrow (E_1 - E_2) \times q_n = 0$$

$$H_{1t} - H_{2t} = K \Rightarrow (H_1 - H_2) \times q_n = K$$

$$D_{1n} - D_{2n} = f_s \Rightarrow (D_1 - D_2) \cdot q_n = f_s$$

$$B_{1n} - B_{2n} = 0 \Rightarrow (B_1 - B_2) \cdot q_n = 0$$



Retarded Potentials

* In the static case we have

$$V = \iiint_V \frac{\rho_V dv}{4\pi\epsilon_0 R}$$

$$\text{and } \vec{A} = \iiint_V \frac{\mu_0 \vec{J} dv}{4\pi R}$$

$$\vec{E} = -\nabla V \quad \text{and} \quad \vec{B} = \nabla \times \vec{A}$$

Retarded Potentials (Cont'd)

* Starting with the E field curl equation, we have

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} = -\frac{\partial}{\partial t} (\nabla \times \vec{A})$$

$$\Rightarrow \nabla \times \left(\vec{E} + \frac{\partial \vec{A}}{\partial t} \right) = 0$$

$$\Rightarrow \vec{E} + \frac{\partial \vec{A}}{\partial t} = -\nabla V \Rightarrow$$

$$\boxed{\vec{E} = -\nabla V - \frac{\partial \vec{A}}{\partial t}}$$

Retarded Potentials (Cont'd)

* The D.E.s governing V and \underline{A} are given by

$$\nabla^2 V + \frac{\partial^2}{\partial t^2} (\nabla \cdot \underline{A}) = -\frac{\rho V}{\epsilon}$$

and

$$\nabla^2 \underline{A} - \nabla (\nabla \cdot \underline{A}) = -\mu \underline{J} + \mu \epsilon \nabla \left(\frac{\partial V}{\partial t} \right) + \mu \epsilon \frac{\partial^2 \underline{A}}{\partial t^2}$$

We choose $\nabla \cdot \underline{A} = -\frac{\partial V}{\partial t}$

Retarded Potential (cont'd)

$$\nabla^2 V - \underline{\underline{\mu\epsilon}} \frac{\partial^2 V}{\partial t^2} = -\frac{\rho_V}{\epsilon}$$

$$\nabla^2 A - \underline{\underline{\mu\epsilon}} \frac{\partial^2 A}{\partial t^2} = -\underline{\underline{\mu J}} \quad \left. \vphantom{\frac{\partial^2 A}{\partial t^2}} \right\} u = \underline{\underline{\frac{1}{\mu\epsilon}}} \text{ m/s}$$

These equations have the solutions

$$V(x, y, z, t) = \iiint_V \frac{\rho_V(\tilde{x}, \tilde{y}, \tilde{z}, t - \frac{R}{u})}{4\pi\epsilon R}$$

$$A(x, y, z, t) = \iiint_V \frac{\mu \underline{\underline{J}}(\tilde{x}, \tilde{y}, \tilde{z}, t - \frac{R}{u})}{4\pi R}$$



Time Harmonic Fields

* If the excitation is sinusoidal, then all fields are sinusoidal quantities

* In this case, we can use phasors

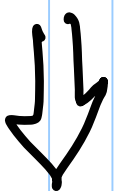
$$E_x = E_0 \cos(\omega t + \phi_x) \\ = \operatorname{Re} (E_0 e^{j(\omega t + \phi_x)})$$

$$E_x = \operatorname{Re} (E_0 e^{j\phi_x}) e^{j\omega t}$$

E_x

Maxwell's Equations For Harmonic Fields

$$\nabla \times \underline{E} = -\frac{\partial \underline{B}}{\partial t}$$



$$\nabla \times \underline{E} = -j\omega \underline{B}$$

$$\nabla \times \underline{H} = \underline{J} + \frac{\partial \underline{D}}{\partial t}$$



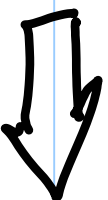
$$\nabla \times \underline{H} = \underline{J} + j\omega \underline{D}$$

$$\nabla \cdot \underline{D} = \rho_v$$



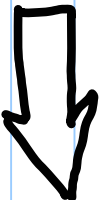
$$\nabla \cdot \underline{D} = \rho_v$$

$$\nabla \cdot \underline{B} = 0$$



$$\nabla \cdot \underline{B} = 0$$

$$\nabla \cdot \underline{J} = -\frac{\partial \rho_v}{\partial t}$$



$$\nabla \cdot \underline{J} = -j\omega \rho_v$$

* Same boundary conditions still apply



EM Power Flow

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$\cdot \vec{H}$ } Subtract

$$\nabla \times \vec{H} = \vec{J}_c + \frac{\partial \vec{D}}{\partial t} \cdot \vec{E}$$

$$\vec{H} \cdot (\nabla \times \vec{E}) - \vec{E} \cdot (\nabla \times \vec{H}) = \vec{H} \cdot \left(-\frac{\partial \vec{B}}{\partial t}\right) - \vec{E} \cdot \left(\vec{J}_c + \frac{\partial \vec{D}}{\partial t}\right)$$

$$\nabla \cdot (\vec{E} \times \vec{H}) = -\vec{H} \cdot \frac{\partial \vec{B}}{\partial t} - \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} - \vec{E} \cdot \vec{J}_c$$

EM Power Flow (Cont'd)

* Typically $\frac{\partial \epsilon}{\partial t} = 0$, $\frac{\partial \mu}{\partial t} = 0$

$$\nabla \cdot (\underline{E} \times \underline{H}) = -\mu \underline{H} \cdot \frac{\partial \underline{H}}{\partial t} - \epsilon \underline{E} \cdot \frac{\partial \underline{E}}{\partial t} - \underline{E} \cdot (\sigma \underline{E})$$

$$\Rightarrow \nabla \cdot (\underline{E} \times \underline{H}) = -\frac{1}{2} \mu \frac{\partial |\underline{H}|^2}{\partial t} - \frac{1}{2} \epsilon \frac{\partial |\underline{E}|^2}{\partial t} - \sigma |\underline{E}|^2$$

$U_e = \frac{1}{2} \epsilon |\underline{E}|^2$ J/m³ electric energy density

$U_m = \frac{1}{2} \mu |\underline{H}|^2$ J/m³ magnetic energy density

$P_d = \sigma |\underline{E}|^2$ W/m³ power density of
Losses

Vector of Poynting

* The vector $\underline{S} = \underline{E} \times \underline{H}$ is called the vector of Poynting

* This vector points in the direction of EM power propagation

* Integrating over the volume of interest, we get

Poynting Theorem

$$\iiint_V \nabla \cdot (\mathbf{E} \times \mathbf{H}) \, dV$$

$$= -\frac{1}{2} \iiint_V \epsilon \frac{\partial |\mathbf{E}|^2}{\partial t} \, dV - \frac{1}{2} \iiint_V \mu \frac{\partial |\mathbf{H}|^2}{\partial t} \, dV$$



$$- \iiint_V \sigma |\mathbf{E}|^2 \, dV$$

$$\begin{aligned} - \iint_S (\mathbf{E} \times \mathbf{H}) \cdot d\mathbf{S} &= \frac{1}{2} \iiint_V \epsilon \frac{\partial |\mathbf{E}|^2}{\partial t} \, dV \\ &+ \frac{1}{2} \iiint_V \mu \frac{\partial |\mathbf{H}|^2}{\partial t} \, dV + \iiint_V \sigma |\mathbf{E}|^2 \, dV \end{aligned}$$

Poynting's Theorem (Cont'd)

$$\begin{aligned} \int_S (\mathbf{E} \times \mathbf{H}) \cdot d\mathbf{S} &= \frac{1}{2} \iiint_V \epsilon \frac{\partial (|\mathbf{E}|^2)}{\partial t} dV \\ &+ \frac{1}{2} \iiint_V m \frac{\partial (|\mathbf{H}|^2)}{\partial t} dV + \iiint_V \sigma |\mathbf{E}|^2 dV \end{aligned}$$

Energy Flowing in through S'

= Stored Energy + dissipated Energy

Poynting's Theorem (Cont'd)

If $\vec{J} = \vec{J}_c + \vec{J}_i$, we get

$$\iiint_V (\vec{E} \cdot \vec{J}_i) dV = \iint_S (\vec{E} \times \vec{H}) \cdot d\vec{S} + \frac{1}{2} \iiint_V \epsilon \frac{\partial (|\vec{E}|^2)}{\partial t} dV + \frac{1}{2} \iiint_V \mu \frac{\partial (|\vec{H}|^2)}{\partial t} dV + \iiint_V \rho (|\vec{E}|^2) dV$$

Energy generated = Energy leaving +

Energy stored + Energy dissipated

EM Energy is Conserved!

