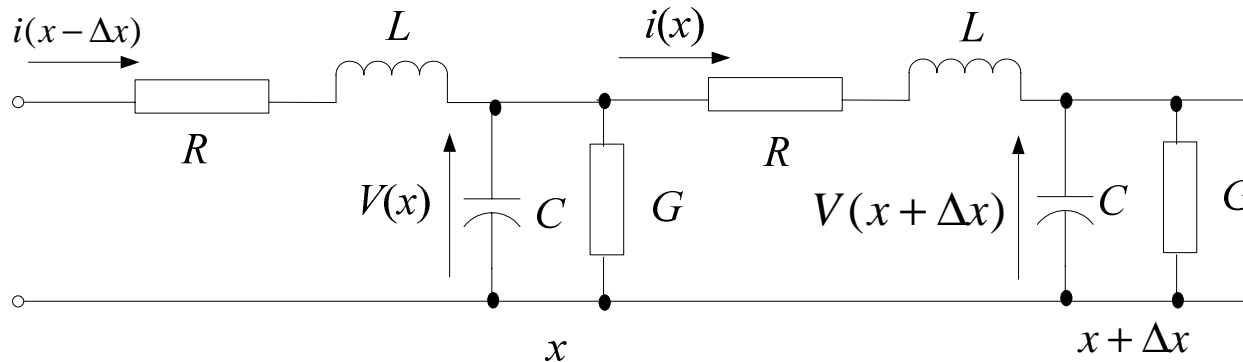


EE750
Advanced Engineering Electromagnetics
Lecture 4

1D TLM

- We establish a one-to-one mapping between 1D wave equations and a network of transmission lines



- C is the capacitance of a section of length Δx , $C = C_d \Delta x$
- L is the inductance of a section of length Δx , $L = L_d \Delta x$
- R and G represent series resistance and shunt conductance, respectively

1D TLM (Cont'd)

- Applying KVL and KCL, we get

$$\Delta x \frac{\partial V}{\partial x} = -L \frac{\partial i}{\partial t} - iR \quad \text{and} \quad \Delta x \frac{\partial i}{\partial x} = -C \frac{\partial V}{\partial t} - GV$$

- Differentiating the first equation w.r.t t and the second equation w.r.t. x we get

$$\frac{\partial^2 i}{\partial x^2} = \frac{GR}{(\Delta x)^2} i + \frac{GL + RC}{(\Delta x)^2} \frac{\partial i}{\partial t} + \frac{LC}{(\Delta x)^2} \frac{\partial^2 i}{\partial t^2}$$

- Similarly, we can show that

$$\frac{\partial^2 V}{\partial x^2} = \frac{GR}{(\Delta x)^2} V + \frac{GL + RC}{(\Delta x)^2} \frac{\partial V}{\partial t} + \frac{LC}{(\Delta x)^2} \frac{\partial^2 V}{\partial t^2}$$

Correspondence with Maxwell's Equations

- For a 1D source-free problem, the fields depend only on one coordinate (say x)

- Maxwell's equations are given by

$$\frac{\partial E_y}{\partial x} = -\frac{\partial B_z}{\partial t} \quad \text{and} \quad \frac{\partial H_z}{\partial x} = J_{cy} + \frac{\partial D_y}{\partial t}$$

- Similarly, we obtain $\frac{\partial^2 E_y}{\partial x^2} = \mu\sigma \frac{\partial E_y}{\partial t} + \mu\epsilon \frac{\partial^2 E_y}{\partial t^2}$

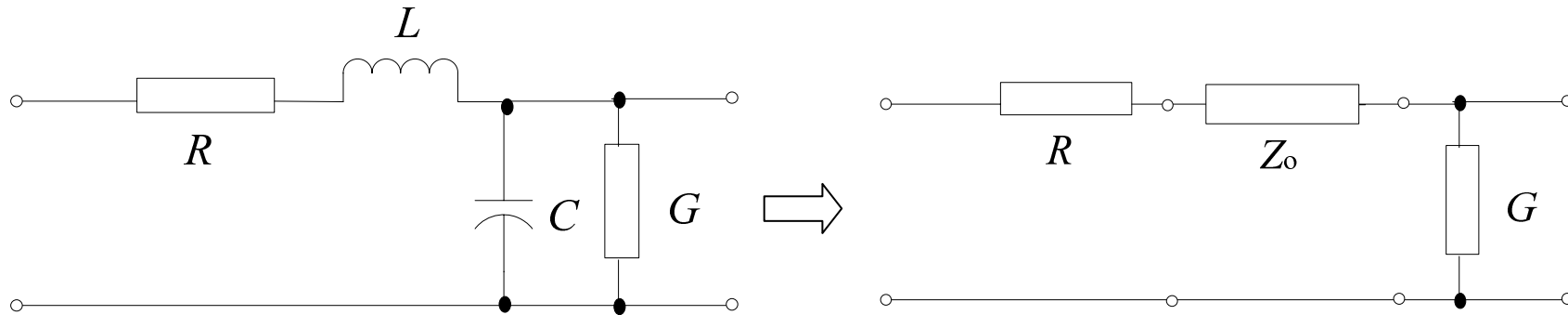
- Comparing with case $R=0$, $\frac{\partial^2 V}{\partial x^2} = \frac{GL}{(\Delta x)^2} \frac{\partial V}{\partial t} + \frac{LC}{(\Delta x)^2} \frac{\partial^2 V}{\partial t^2}$

We obtain the one-to-one correspondence

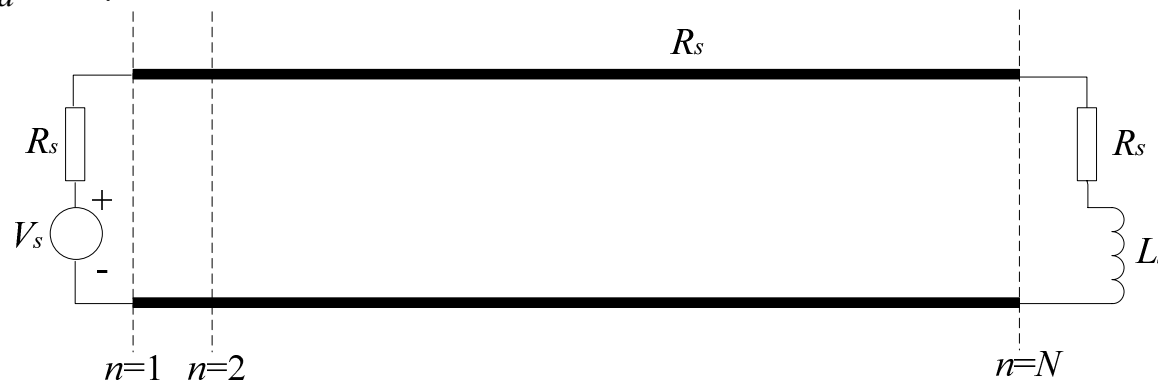
$$V \leftrightarrow E_y, \quad \mu \leftrightarrow (L/\Delta x), \quad \epsilon \leftrightarrow (C/\Delta x), \quad \sigma \leftrightarrow (G/\Delta x)$$

- Solving the discretized TLM network obtains a solution of the corresponding EM problem

Solution of the TLM Network



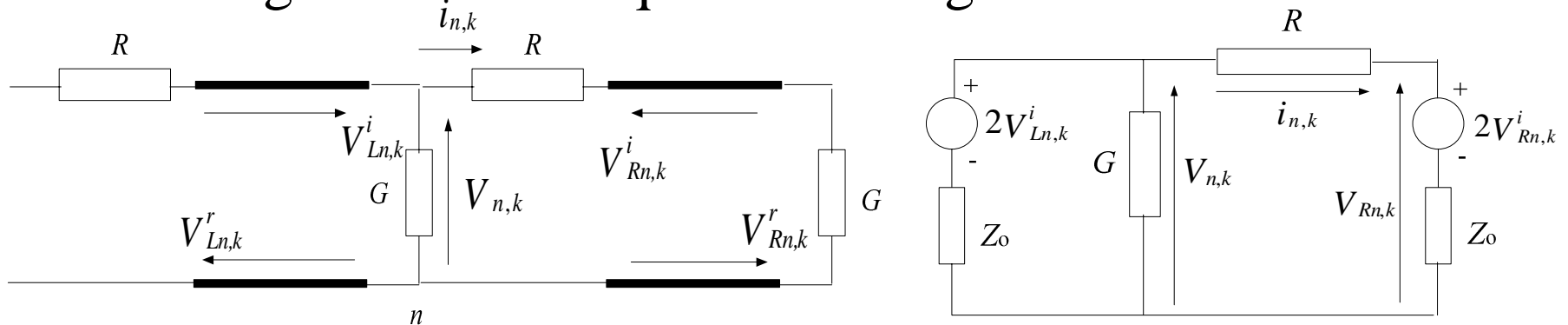
$$Z_o = \sqrt{\frac{L_d}{C_d}} = \sqrt{\frac{L}{C}}, \quad \Delta t = \frac{\Delta x}{u} = \Delta x \sqrt{L_d C_d} = \sqrt{LC}$$



- Consider the following model with N sections
- Node n is between sections $(n-1)$ and n , $1 < n < (N+1)$

Solution of a TLM Network (Cont'd)

- Utilizing Thevenin's equivalent we get



- Using Superposition or Milliman's theorem, we get

$$V_{n,k} = \frac{\frac{2V_{Ln,k}^i}{Z_o} + \frac{2V_{Rn,k}^i}{Z_o}}{\frac{1}{Z_o} + \frac{1}{G} + \frac{1}{Z_o + R}} \quad \Rightarrow \quad i_{n,k} = \frac{V_{n,k} - 2V_{Rn,k}^i}{Z_o + R}$$

$$V_{Rn,k} = 2V_{Rn,k}^i + i_{n,k} Z_o \quad , \quad V_{Ln,k} = V_{n,k}$$

Solution of a TLM Network (Cont'd)

- Using all the voltages and current, we carry out the two fundamental steps of the TLM method:
- Scattering: Evaluate the reflected waves

$$V_{Rn,k}^r = V_{Rn,k} - V_{Rn,k}^i, \quad V_{Ln,k}^r = V_{Ln,k} - V_{Ln,k}^i$$

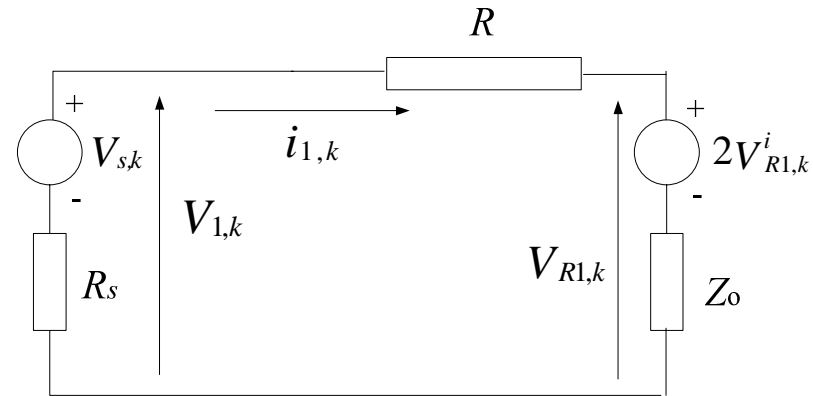
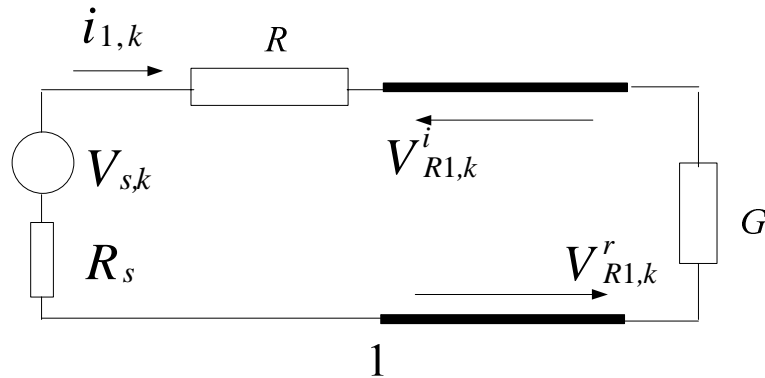
- Connection: determine the incident waves at the $(k+1)$ time step \Rightarrow reflected waves become incident waves on neighboring nodes at the next time step

$$V_{Ln,(k+1)}^i = V_{R(n-1),k}^r$$

$$V_{Rn,(k+1)}^i = V_{L(n+1),k}^r$$

Solution of a TLM Network (Cont'd)

- For the source node we have

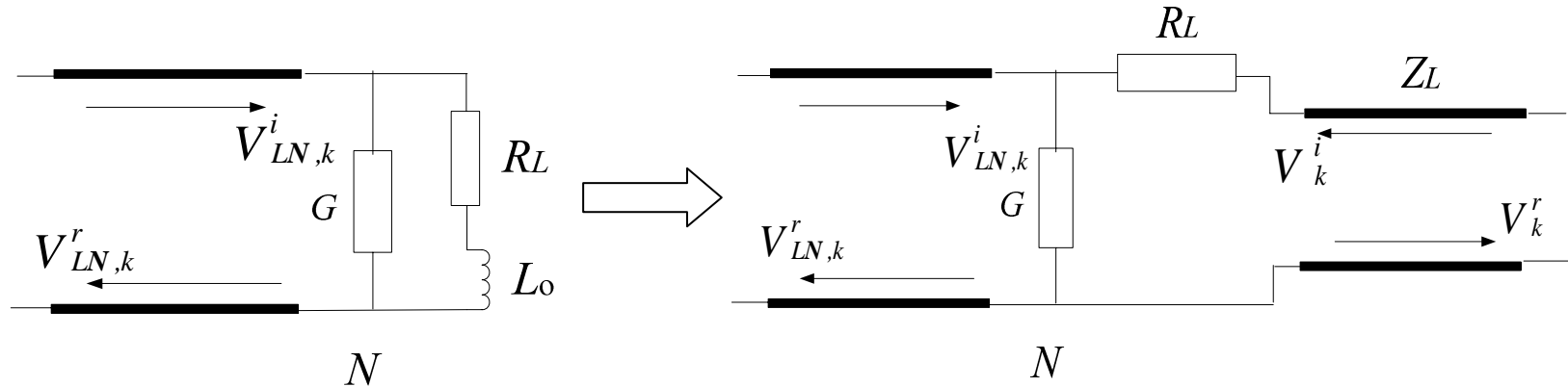


$$i_{1,k} = \frac{V_{s,k} - 2V_{R1,k}^i}{R + R_s + Z_0}, \quad V_{R1,k} = 2V_{R1,k}^i + i_{1,k} Z_0$$

$$V_{R1,k}^r = V_{R1,k} - V_{R1,k}^i, \quad V_{R1,k+1}^i = V_{L2,k}^r$$

Solution of the TLM Network (Cont'd)

- For the load node we have



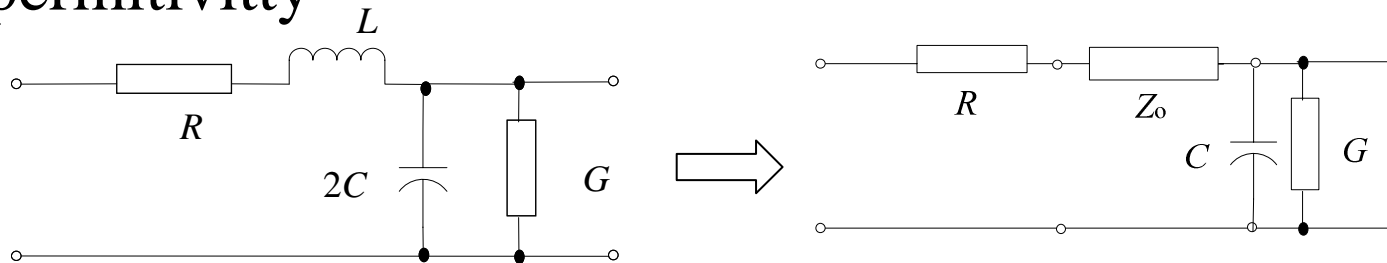
- Note that $Z_L = L_o / (\Delta t / 2) \implies$ synchronization is preserved
- Derive the scattering and connection relationships for the load node

General Steps

- Given the parameters of the electromagnetic problem determine the parameters of the TLM network (Source, load, time step, L , C)
- Initialize incident impulses for all sections (left and right)
- Repeat for all time steps
 - * Evaluate intermediate quantities (voltages and currents)
 - * Evaluate reflected waves (scattering)
 - * Obtain incident waves at the next time step from the reflected impulses at the current time step (Connection)end

The Nonhomogenous Case

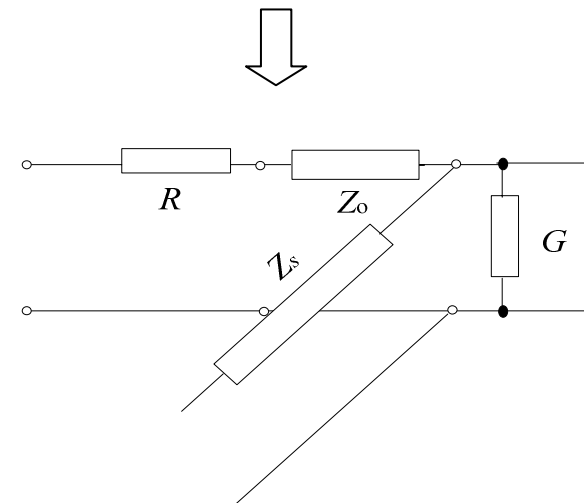
- To maintain synchronization, stubs should be used
- Example: a section of length Δx with double the permittivity



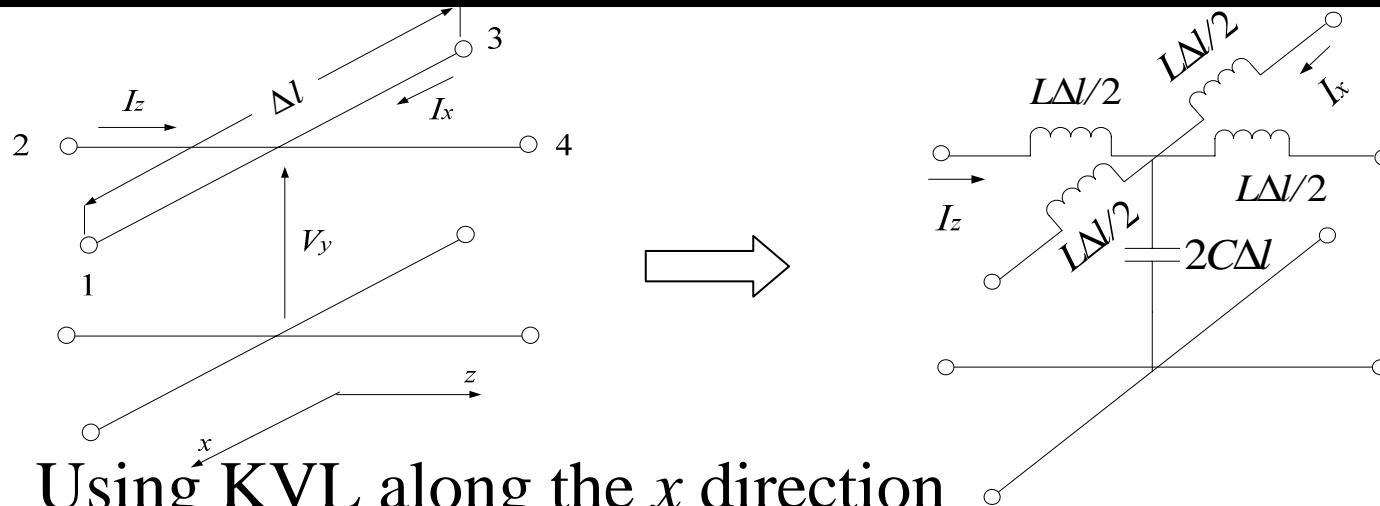
$$Z_0 = \sqrt{\frac{L}{C}} \quad , \quad u = \frac{\Delta x}{\sqrt{LC}}$$

$$\Delta t = \frac{\Delta x}{u} = \sqrt{LC}$$

For the stub we have $Z_s = \frac{\Delta t}{C}$



2D TLM using Shunt Nodes



- Using KVL along the x direction

$$V_y(x + \Delta x) = V_y(x) - L\Delta x \frac{\partial I_x}{\partial t} \implies \boxed{\frac{\partial V_y}{\partial x} = -L \frac{\partial I_x}{\partial t}}$$

- Using KVL along the z direction we have

$$V_y(z + \Delta z) = V_y(z) - L\Delta z \frac{\partial I_z}{\partial t} \implies \boxed{\frac{\partial V_y}{\partial z} = -L \frac{\partial I_z}{\partial t}}$$

- Using KCL we have $I_z(z) + I_x(x) = I_z(z + \Delta z) + I_x(x + \Delta x) + 2C\Delta l \frac{\partial V_y}{\partial t}$

$$\implies -\frac{\partial I_z}{\partial z} \Delta l - \frac{\partial I_x}{\partial x} \Delta l = 2C\Delta l \frac{\partial V_y}{\partial t} \implies \boxed{\frac{\partial I_z}{\partial z} + \frac{\partial I_x}{\partial x} = -2C \frac{\partial V_y}{\partial t}}$$

2D TLM using Shunt Nodes (Cont'd)

- Combining these equations we get $\frac{\partial^2 V_y}{\partial x^2} + \frac{\partial^2 V_y}{\partial z^2} = 2LC \frac{\partial^2 V_y}{\partial t^2}$
- Maxwell's equations for the lossless TM_y case are

$$\frac{\partial E_y}{\partial z} = \mu \frac{\partial H_x}{\partial t}, \quad \frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} = \epsilon \frac{\partial E_y}{\partial t}, \quad \frac{\partial E_y}{\partial x} = -\mu \frac{\partial H_z}{\partial t}$$
- Combining these equations, we get $\frac{\partial^2 E_y}{\partial x^2} + \frac{\partial^2 E_y}{\partial z^2} = 2LC \frac{\partial^2 E_y}{\partial t^2}$
- We now establish the equivalences

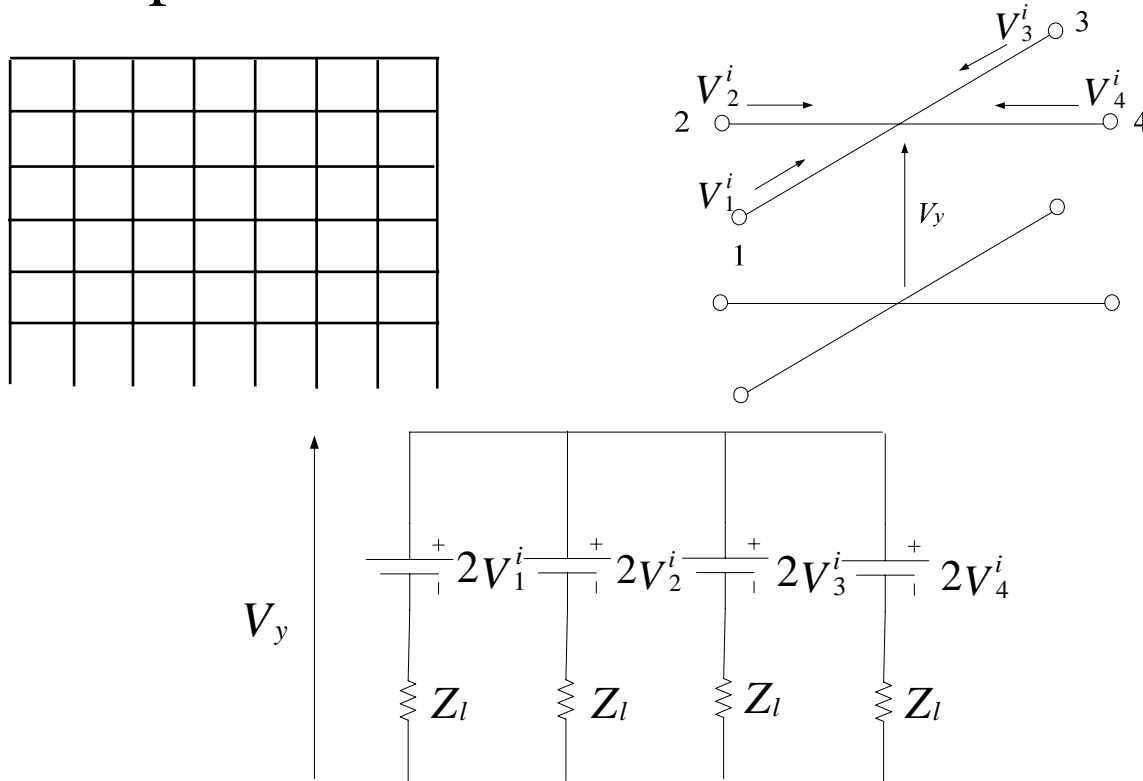
$$E_y \leftrightarrow V_y, \quad H_z \leftrightarrow I_x, \quad H_x \leftrightarrow -I_z, \quad \mu \leftrightarrow L, \quad \epsilon \leftrightarrow 2C$$

- network velocity=medium velocity $1/\sqrt{\mu\epsilon} = 1/\sqrt{2CL}$
- The link velocity v_l is $v_l = 1/\sqrt{LC} = \sqrt{2}/\sqrt{\mu\epsilon} = \sqrt{2} v_n$
- The link characteristic impedance is

$$Z_l = \sqrt{L/C} = \sqrt{2}\sqrt{\mu/\epsilon} = \sqrt{2} Z_n$$

Scattering in Shunt TLM Nodes (Lossless Case)

- The computational domain is filled with TLM cells



- We replace each section by its Thevenin's equivalent
- Using Superposition we have $V_y = 0.5(V_1^i + V_2^i + V_3^i + V_4^i)$

Scattering in Shunt TLM Nodes (Cont'd)

- It follows that we have

$$V_1^r = V_y - V_1^i = 0.5(-V_1^i + V_2^i + V_3^i + V_4^i)$$

$$V_2^r = V_y - V_2^i = 0.5(V_1^i - V_2^i + V_3^i + V_4^i)$$

$$V_3^r = V_y - V_3^i = 0.5(V_1^i + V_2^i - V_3^i + V_4^i)$$

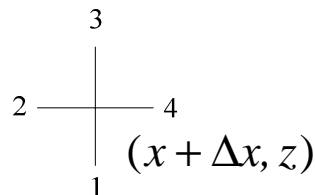
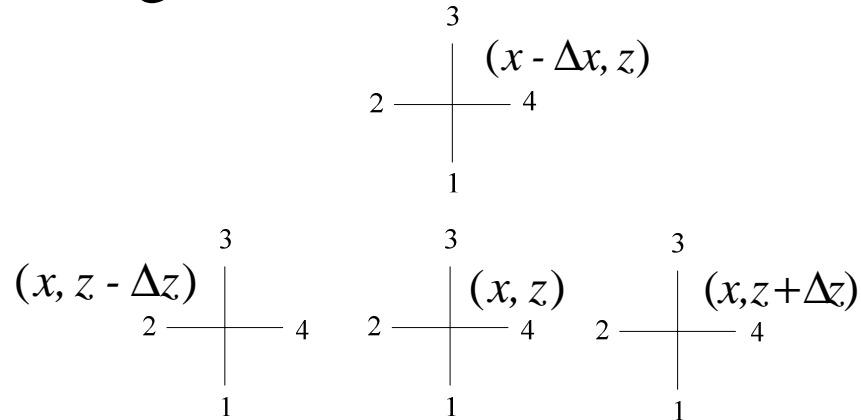
$$V_4^r = V_y - V_4^i = 0.5(V_1^i + V_2^i + V_3^i - V_4^i)$$

- Or in matrix form

$$\begin{bmatrix} V_1^r \\ V_2^r \\ V_3^r \\ V_4^r \end{bmatrix} = 0.5 \begin{bmatrix} -1 & 1 & 1 & 1 \\ 1 & -1 & 1 & 1 \\ 1 & 1 & -1 & 1 \\ 1 & 1 & 1 & -1 \end{bmatrix} \begin{bmatrix} V_1^i \\ V_2^i \\ V_3^i \\ V_4^i \end{bmatrix}$$

Connection in Shunt TLM Nodes

- Reflected impulses at the k th time step become incident on neighboring nodes at the $k+1$ time step



$$\left(V_1^i(x, z)\right)_{k+1} = \left(V_3^r(x + \Delta x, z)\right)_k, \quad \left(V_2^i(x, z)\right)_{k+1} = \left(V_4^r(x, z - \Delta z)\right)_k$$

$$\left(V_3^i(x, z)\right)_{k+1} = \left(V_1^r(x - \Delta x, z)\right)_k, \quad \left(V_4^i(x, z)\right)_{k+1} = \left(V_2^r(x, z + \Delta z)\right)_k$$

Modeling of Free Space

- To model free space, the following two conditions must be satisfied

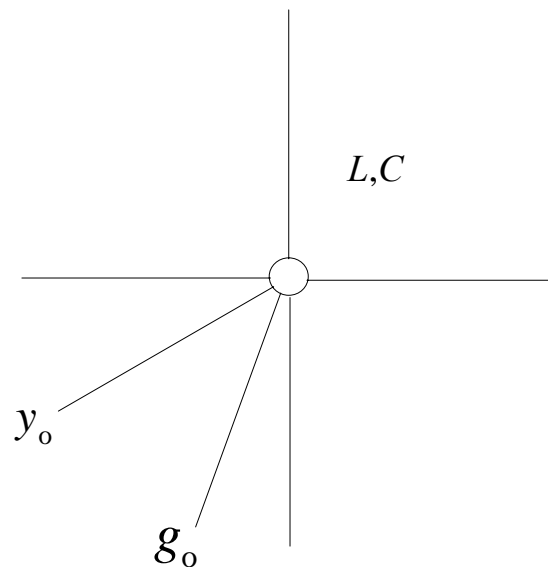
$$\frac{1}{\sqrt{2LC}} = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \implies \sqrt{2LC} = \frac{1}{v_0} = 3.0 \times 10^8 \text{ m/s}$$

$$\sqrt{\frac{L}{2C}} = \sqrt{\frac{\mu_0}{\epsilon_0}} = \eta_0 = 377 \Omega$$

- It follows that $L = \eta_0 / v_0 = \mu_0$ and $C = 1/(2\eta_0 v_0) = \epsilon_0 / 2$
- $v_L = \text{link velocity} = 1/\sqrt{LC} = v_0 \sqrt{2} \text{ m/s}$
- $Z_L = \text{link characteristic impedance} = \sqrt{L/C} = \sqrt{2} \eta_0 \Omega$
- Unit time step $\Delta t = \Delta l / v_L = \Delta l / (v_0 \sqrt{2}) \text{ sec}$

Modeling of a General Lossy Medium

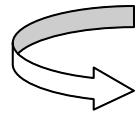
- Additional losses and Capacitances are modeled by matched shunt stubs and open ended shunt stubs, respectively



- y_o : normalized admittance of the $\Delta l/2$ shunt capacitance stub
- g_o : normalized admittance of the $\Delta l/2$ shunt capacitance stub

Modeling of a General Lossy Medium (Cont'd)

- Normalizing admittance is $Y_L = \sqrt{C/L}$



$$Y_o = y_o \sqrt{C/L} \quad \text{and} \quad G_o = g_o \sqrt{C/L}$$

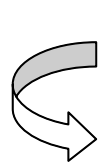
- We choose C_s and L_s for the permittivity stub to be

$$C_s = Cy_o \quad \text{and} \quad L_s = L/y_o \quad \implies \quad 1/\sqrt{L_s C_s} = 1/\sqrt{LC} \quad , \quad \sqrt{\frac{C_s}{L_s}} = Y_o = y_o Y_L$$

- Similarly for the loss stub we have

$$C_g = Cg_o \quad \text{and} \quad L_g = L/g_o \quad \implies \quad 1/\sqrt{L_g C_g} = 1/\sqrt{LC} \quad , \quad \sqrt{\frac{C_g}{L_g}} = G_o = g_o Y_L$$

- The permittivity stub represents a capacitance of value

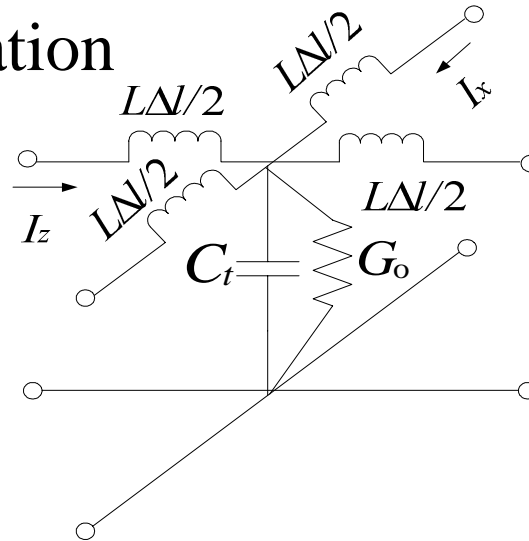


$$C_s \Delta l/2 = Cy_o \Delta l/2 \quad \implies \quad \text{Total cell capacitance} = Cy_o \Delta l/2 + 2C$$

$$C_t = 2C\Delta l \left(1 + \frac{y_o}{4}\right)$$

Modeling of a General Lossy Medium (Cont'd)

- For small cell size, the cell now has the equivalent lumped presentation



- Equations governing this lossy node are

$$\frac{\partial V_y}{\partial x} = -L \frac{\partial I_x}{\partial t} \quad (\text{KVL in the } x \text{ direction})$$

$$\frac{\partial V_y}{\partial z} = -L \frac{\partial I_z}{\partial t} \quad (\text{KVL in the } z \text{ direction})$$

$$\frac{\partial I_z}{\partial z} + \frac{\partial I_x}{\partial x} = -2C(1 + y_o/4) \frac{\partial V_y}{\partial t} - \frac{g_o \sqrt{C/L}}{\Delta l} V_y \quad (\text{KCL})$$

Correspondence with Maxwell's Equations

- Maxwell's equations for the lossy 2D TE_y case are

$$\partial E_y / \partial x = -\mu \partial H_z / \partial t$$

$$\partial E_y / \partial z = \mu \partial H_x / \partial t$$

$$\frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} = \varepsilon \frac{\partial E_y}{\partial t} + \sigma E_y$$

- It follows that we can establish the 1-1 correspondence

$$E_y \leftrightarrow V_y, \quad H_z \leftrightarrow I_x, \quad H_x \leftrightarrow -I_z$$

$$\mu \leftrightarrow L, \quad \varepsilon \leftrightarrow 2C\left(1 + \frac{y_0}{4}\right), \quad \sigma \leftrightarrow \frac{g_0 \sqrt{C/L}}{\Delta l}$$

Some Practical Points

- L and C are usually selected to model free space with $L=\mu_0$ and $C=\epsilon_0/2$
- y_0 is adjusted at each node to model the local permittivity
$$\epsilon = 2C\left(1 + \frac{y_0}{4}\right) \implies \epsilon_r = \left(1 + \frac{y_0}{4}\right) \implies y_0 = 4.0(\epsilon_r - 1)$$
- g_0 is adjusted at each node to model the local conductivity

$$\sigma = \frac{g_0 \sqrt{C/L}}{\Delta l}$$